# The Spectral Gap of the 2-D Stochastic Ising Model with Nearly Single-Spin Boundary Conditions 

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#### Abstract

We establish upper bounds for the spectral gap of the stochastic Ising model at low temperature in an $N \times N$ box, with boundary conditions which are "plus" except for small regions at the corners which are either free or "minus." The spectral gap decreases exponentially in the size of the corner regions, when these regions are of size at least of order $\log N$. This means that removing as few as $O(\log N)$ plus spins from the corners produces a spectral gap far smaller than the order $N^{-2}$ gap believed to hold under the all-plus boundary condition. Our results are valid at all subcritical temperatures.


KEY WORDS: Stochastic Ising model; spectral gap; Glauber dynamics.

## 1. INTRODUCTION AND MAIN THEOREM

Let $\Lambda \subset \mathbb{Z}^{2}$ and let $\eta \in\{-1,0,1\}^{\partial \Lambda}$. Here $\partial \Lambda=\left\{x \in \mathbb{Z}^{2} \backslash \Lambda: x\right.$ is adjacent to some site in $\Lambda\}$. The Hamiltonian for the Ising model on $\Lambda$ with boundary condition $\eta$ is

$$
H_{A, \eta}(\sigma)=-\sum_{\langle x\rangle\rangle: x, y \in \Lambda} \sigma_{x} \sigma_{y}-\sum_{\langle x y\rangle: x \in \Lambda, y \in \partial \Lambda} \sigma_{x} \eta_{y}, \quad \sigma \in\{-1,1\}^{\Lambda},
$$

where the first sum is over unordered pairs of adjacent sites. Let $\mu=\mu_{\Lambda, \eta}^{\beta}$ denote the equilibrium measure when the inverse temperature is $\beta$ :

$$
\mu_{A, \eta}^{\beta}(\sigma)=\left(Z_{\Lambda, \eta}^{\beta}\right)^{-1} e^{-\beta H_{\Lambda, \eta}(\sigma)},
$$

[^0]where $Z_{A, \eta}^{\beta}$ is the partition function. Let $\Sigma_{\Lambda}=\{-1,1\}^{4}$ be the configuration space, and let $\sigma=\sigma_{A}$ denote a generic configuration. (When convenient we write " + " and " - " in place of 1 and -1 for the two spins.)

We consider the time evolution of the dynamic version of the model under Glauber dynamics. Let $\sigma^{x}$ denote the configuration given by $\sigma_{y}^{x}=-\sigma_{y}$ for $y=x, \sigma_{y}^{x}=\sigma_{y}$ for $y \neq x$. The flip rate at a site $x$ when the configuration is $\sigma$ is denoted $c(x, \sigma)$ (notationally supressing its possible dependence on $\Lambda, \eta$.) We assume that the flip rates are uniformly bounded:

$$
0<c_{0}^{\prime} \leqslant c(x, \sigma) \leqslant c_{0} \quad \text { for all } \quad x, \sigma, \Lambda, \eta
$$

We also make the usual assumptions that the flip rates are attractive and translation invariant, satisfy detailed balance and have finite range; see, e.g., ref. 15 for full descriptions of these properties. The generator $A=A_{1, \eta}^{\beta}$ of the corresponding Markov process is given by

$$
(A f)(\sigma)=\sum_{x \in \Lambda} c(x, \sigma)\left(f\left(\sigma^{x}\right)-f(\sigma)\right)
$$

and the Dirichlet form $\mathscr{D}=\mathscr{D}_{1, \eta}^{\beta}$ by

$$
\mathscr{D}(f, g)=\langle f, A g\rangle_{\mu},
$$

so that

$$
\mathscr{D}(f, f)=-\frac{1}{2} \sum_{x \in \Lambda} \sum_{\sigma \in \Sigma_{A}} \mu(\sigma) c(x, \sigma)\left(f\left(\sigma^{x}\right)-f(\sigma)\right)^{2},
$$

where $\mu=\mu_{\Lambda, \eta}^{\beta}$. The spectral gap $\Delta=\Delta(\Lambda, \eta, \beta)$ for such dynamics, which is the smallest positive eigenvalue of $-A_{\Lambda, \eta}^{\beta}$, has the following representation:

$$
\begin{equation*}
\Delta(\Lambda, \eta, \beta)=\inf _{f \in L^{2}(\mu)}-\frac{\mathscr{D}(f, f)}{\operatorname{var}_{\mu}(f)}, \tag{1.1}
\end{equation*}
$$

where $\operatorname{var}_{\mu}(f)$ denotes the variance of $f$. The gap $\Delta$ describes the rate of exponential convergence in $L^{2}\left(\mu_{\Lambda, \eta}^{\beta}\right)$ to equilibrium, in the sense that for $S(\cdot)$ the semigroup generated by $A$ and $\|\cdot\|_{\mu}$ the $L^{2}(\mu)$ norm, $\Delta$ is the largest constant such that

$$
\left\|S(t) f-\int f d \mu\right\|_{\mu} \leqslant\left\|f-\int f d \mu\right\|_{\mu} e^{-\Delta t} \quad \text { for all } \quad f \in L^{2}(\mu) \text { and } t \geqslant 0 .
$$

We say that two configurations $\sigma, \sigma^{\prime} \in \Sigma_{A}$ are adjacent if for some $x \in \Lambda$ we have $\sigma_{x} \neq \sigma_{x}^{\prime}$ but $\sigma_{y}=\sigma_{y}^{\prime}$ for all $y \neq x$. For $S \subset \Sigma_{A}$ define

$$
\partial_{i n} S=\left\{\sigma \in S: \sigma \text { is adjacent to some configuration in } S^{c}\right\} .
$$

Considering only indicator functions we obtain

$$
\begin{equation*}
\Delta(\Lambda, \eta, \beta) \leqslant c_{0}|\Lambda| \inf _{D \subset \Sigma_{\Lambda}} \frac{\mu_{\Lambda, \eta}^{\beta}\left(\partial_{i n} D\right)}{\mu_{\Lambda, \eta}^{\beta}(D)\left(1-\mu_{\Lambda, \eta}^{\beta}(D)\right)} \tag{1.2}
\end{equation*}
$$

Let $\tilde{\Lambda}_{N}=[-N, N]^{2}$ and $\Lambda_{N}=\tilde{\Lambda}_{N} \cap \mathbb{Z}^{2}$. Let $l_{h}$ and $l_{v}$ denote the horizontal and vertical axes, respectively, and consider the boundary condition $\eta^{k, \epsilon}$ given for $k \geqslant 0$ and $\epsilon \in\{0,-1\}$ by

$$
\eta_{x}^{k, \epsilon}= \begin{cases}1, & \text { if } d\left(x, l_{h}\right) \leqslant k \quad \text { or } \quad d\left(x, l_{v}\right) \leqslant k ;  \tag{1.3}\\ \epsilon, & \text { otherwise } .\end{cases}
$$

Here $d(\cdot, \cdot)$ denotes Euclidean distance. As a special case of results in ref. 13 we have that for $\beta$ very large, for some $C, \lambda$ depending only on $\delta, \beta$ and $\epsilon$,

$$
\begin{equation*}
\Delta\left(\Lambda_{N}, \eta^{\delta N, \epsilon}, \beta\right) \leqslant C e^{-\lambda N} \quad \text { for all } \quad \delta \leqslant \frac{1}{8} \quad \text { and } \quad N \geqslant 1 . \tag{1.4}
\end{equation*}
$$

Here we will generalize this as follows. Let $\beta_{c}$ denote the critical inverse temperature of the Ising model on $\mathbb{Z}^{2}$.

Theorem 1.1. Let $\beta>\beta_{c}$.
(i) For some $C, K, \lambda$ depending only on $\beta$, for all $N \geqslant 1$ and $k \geqslant 1$ satisfying $N-k \geqslant K \log N$,

$$
\begin{equation*}
\Delta\left(\Lambda_{N}, \eta^{k, 0}, \beta\right) \leqslant C e^{-\lambda(N-k)} . \tag{1.5}
\end{equation*}
$$

(ii) For some $C, K, \lambda$ depending only on $\beta$, for all $N \geqslant 1$ and $k \geqslant 1$ satisfying $\min (k, N-k) \geqslant K \log N$,

$$
\begin{equation*}
\Delta\left(\Lambda_{N}, \eta^{k,-1}, \beta\right) \leqslant C e^{-\lambda \min (k, N-k)} . \tag{1.6}
\end{equation*}
$$

By Theorem 1.2 below, the spectral gap changes by at most a constant when a single boundary spin is changed; in particular this can be applied in comparing $k=1$ to a completely free boundary. Thus Theorem 1.1(i) essentially includes exponential decay of the gap under the free boundary condition at low temperatures, a result obtained by Thomas. ${ }^{(20)}$

Theorem 1.1 generalizes (1.4) to all $\delta<1$ and $\beta>\beta_{c}$, and shows that the $L^{2}$ rate of convergence to equilibrium in the stochastic Ising model can be quite slow even when the boundary condition is overwhelmingly of a single spin, and otherwise free. For the full "plus" boundary condition ( $k=N$ ) at subcritical temperatures, it is believed ${ }^{(7)}$ that the spectral gap is of order $N^{-2}$; Martinelli ${ }^{(16)}$ proved that for very low temperatures,

$$
\begin{equation*}
\Delta\left(\Lambda_{N},+, \beta\right) \geqslant \exp \left(-\lambda(\delta) N^{\frac{1}{2}+\delta}\right) \quad \text { for all } \quad \delta>0 . \tag{1.7}
\end{equation*}
$$

Presuming $N^{-2}$ is the correct rate for the "plus" boundary condition, Theorem 1.1(i) shows that removing as few as $O(\log N)$ plus spins from the corners of the box dramatically shrinks the spectral gap.

For $k \approx c N$ with $0<c<1$, Theorem 1.1 shows that the spectral gap decreases at least exponentially fast in $N$. Schonmann ${ }^{(19)}$ showed that for all $\eta$, all $N$ and all $\beta>0$, for some $C=C(\beta)$,

$$
\Delta\left(\Lambda_{N}, \eta, \beta\right) \geqslant \frac{C}{N} e^{-4 \beta N}
$$

so the gap can never decrease faster than exponentially in $N$.
Our proof of Theorem 1.1 will use the method suggested by (1.2): we find an event $D$ with $\partial_{\text {in }} D$ much smaller than $D$. This event $D$ is a variant of the event that none of the four strips of " + " spins in $\partial \Lambda_{N}$ is connected to any of the other strips by a path of " + " spins.

It would be of interest to establish similar lower bounds to go with the upper bounds of Theorem 1.1, but this would most likely entail solving the difficult open problem of obtaining a good lower bound for the gap for full "plus" boundary conditions $(k=N)$. With (1.7) now the best available result and $N^{-2}$ the conjectured actual decay rate, such a bound does not appear obtainable at present. However, the following simple result enables us to relate $k=N$ to smaller values of $k$, and shows conditionally on the validity of the conjectured $N^{-2}$ rate that lower bounds complementing (1.5) and (1.6) do hold.

Theorem 1.2. Let $\beta>0$.
(i) There exists $R=R(\beta)>0$ such that for all $N$, all $\eta \in\{-1,0,1\}^{\Lambda_{N}}$ and all $x \in \partial \Lambda_{N}$, for $\eta^{\prime}$ satisfying $\eta_{y}^{\prime}=\eta_{y}$ for all $y \neq x$ in $\partial \Lambda_{N}$,

$$
\Delta\left(\Lambda_{N}, \eta^{\prime}, \beta\right) \leqslant R \Delta\left(\Lambda_{N}, \eta, \beta\right) .
$$

(ii) For some $C, \lambda$ depending only on $\beta$, for all $1 \leqslant k \leqslant N$,

$$
\begin{aligned}
& \Delta\left(\Lambda_{N}, \eta^{k, 0}, \beta\right) \geqslant C e^{-\lambda(N-k)} \Delta\left(\Lambda_{N},+, \beta\right), \\
& \Delta\left(\Lambda_{N}, \eta^{k,-1}, \beta\right) \geqslant C e^{-\lambda \min (k, N-k)} \Delta\left(\Lambda_{N},+, \beta\right),
\end{aligned}
$$

Here " + " denotes the all-plus boundary condition.

## 2. PRELIMINARIES

Throughout the paper, $c_{0}, c_{1}, \ldots$ and $\epsilon_{1}, \epsilon_{2}, \ldots$ will be used to represent constants which depend only on the temperature (or other parameters) of the model. We use $\epsilon_{i}$ for constants which should be viewed as "small."

Our proof will make use of the Fortuin-Kasteleyn random cluster model, or briefly, the $F K$ model (refs. $8-10$; see also refs. 1 and 11) which is a graphical representation of the Potts model. To discuss this, we need some notation for bond configurations. By a bond we mean an unordered pair $\langle x y\rangle$ of adjacent sites in $\mathbb{Z}^{2}$. When convenient we view bonds as being open line segments in the plane; this should be clear from the context. Define the sets of bonds

$$
\mathscr{B}(\Lambda)=\{\langle x y\rangle: x, y \in \Lambda\}, \quad \overline{\mathscr{B}}(\Lambda)=\{\langle x y\rangle: x \in \Lambda \text { or } y \in \Lambda\} .
$$

For general $A \subset \mathbb{R}^{2}$, we write $\mathscr{B}(A)$ for $\mathscr{B}\left(A \cap \mathbb{Z}^{2}\right)$. Let $\Omega_{\Lambda}=\{0,1\}^{\overline{\mathscr{B}}(A)}$. A bond configuration is an element $\omega \in \Omega_{1}$; when convenient we alternatively view $\omega$ as a subset of $\overline{\mathscr{B}}(\Lambda)$ or as a subgraph of $(\bar{\Lambda}, \overline{\mathscr{B}}(\Lambda))$. Bonds $e$ with $\omega_{e}=1$ are open in $\omega$; those with $\omega_{e}=0$ are closed. Let $C(\omega)$ denote the number of open clusters in $\omega$ which do not intersect $\partial \Lambda$. For $p \in[0,1]$ and $q>0$, the FK model $P_{\Lambda, w}^{p, q}$ on $(\bar{\Lambda}, \overline{\mathscr{B}}(\Lambda))$ with parameters $(p, q)$ and wired boundary condition is defined by the weights

$$
\begin{equation*}
W(\omega)=p^{|\omega|}(1-p)^{|\overline{\mathscr{F}}(1)|-|\omega|} q^{C(\omega)} \tag{2.1}
\end{equation*}
$$

Here $|\omega|$ means the number of open bonds in $\omega$. More generally, given $\rho \in\{0,1\}^{\overline{\operatorname{s}}(1)^{c}}$ we define $(\omega \rho)$ to be the bond configuration on the full lattice which coincides with $\omega$ on $\overline{\mathscr{B}}(\Lambda)$ and with $\rho$ on $\overline{\mathscr{B}}(\Lambda)^{c}$. Let $C(\omega \mid \rho)$ be the number of open clusters of $(\omega \rho)$ which intersect $\Lambda$. The FK model $P_{\Lambda, \rho}^{p, q}$ with bond boundary condition $\rho$ is given by the weights in (2.1) with $C(\omega)$ replaced by $C(\omega \mid \rho)$. Alternately, given $\eta \in\{-1,0,1\}^{\partial \Lambda}$ define

$$
\begin{align*}
V(\Lambda, \eta)= & \left\{\omega \in\{0,1\}^{\bar{s}(1)}: \eta_{x}=\eta_{y} \text { for every } x, y \in \partial \Lambda \text { for which } x \leftrightarrow y \text { in } \omega,\right. \\
& \left.\omega_{e}=0 \text { for all } e \in\left\{\langle x y\rangle: x \in \Lambda, y \in \partial \Lambda, \eta_{y}=0\right\}\right\} . \tag{2.2}
\end{align*}
$$

Here $x \leftrightarrow y$ means there is a path of open bonds connecting $x$ to $y$. The FK model $P_{i, \eta}^{p, q}$ with site boundary condition $\eta$ is given by the weights in (2.1), multiplied by $\delta_{V(1, \eta)}(\omega)$. Taking $\eta_{x}=0$ for all $x$ gives the FK model with free boundary condition; we denote it $P_{A, f}^{p, q}$. For a summary of basic properties of the FK model, see ref. 11. In particular, since we are in two dimensions, for $p \neq \sqrt{q} /(1+\sqrt{q})$ there is a unique translation-invariant infinite-volume FK measure on $\mathscr{B}\left(\mathbb{Z}^{2}\right)$, which can be obtained as the limit of $P_{\Lambda, w}^{p, q}$ as $\Lambda \nearrow \mathbb{Z}^{2}$; we denote this measure $P^{p, q}$. For $q \geqslant 1$, the FK model has the FKG property. For

$$
\alpha(p, q)=\frac{p}{p+q(1-p)}
$$

and $e \in \mathscr{B}\left(\mathbb{Z}^{2}\right)$, we have

$$
P^{p, q}\left(\omega_{e}=1 \mid \omega_{b}, b \neq e\right) \geqslant \alpha(p, q) \quad \text { for every }\left(\omega_{b}, b \neq e\right) .
$$

Changing a single bond in the boundary condition changes the value of $C(\omega \mid \rho)$ by at most 1 . It follows easily that for boundary conditions $\rho, \rho^{\prime}$ differing at only one bond, we have

$$
\begin{equation*}
P_{A, p}^{p, q} \leqslant q P_{A, p^{\prime}}^{p, q} . \tag{2.3}
\end{equation*}
$$

As shown in ref. 6, for $\beta$ given by $p=1-e^{-\beta}$, a configuration of the Ising model on $\Lambda$ with boundary condition $\eta$ at inverse temperature $\beta$ can be obtained from a configuration $\omega$ of the FK model at $(p, 2)$ with site boundary condition $\eta$, by choosing a label for each cluster of $\omega$ independently and uniformly from $\{-1,1\}$; this cluster-labeling construction yields a joint site-bond configuration for which the sites are an Ising model and the bonds are an FK model. When the parameters are related in this way, we call the Ising and FK models corresponding. Alternately, if one selects an Ising configuration $\sigma$ and does independent percolation at density $p$ on the set of bonds

$$
\left\{\langle x y\rangle \in \overline{\mathscr{B}}(\Lambda): \sigma_{x}=\sigma_{y}\right\},
$$

the resulting bond configuration is a realization of the corresponding FK model. We call this the percolation construction of the FK model. From

Onsager's exact solution of the Ising model (see ref. 18) and basic properties of the FK model (see ref. 11) the critical point $\beta_{c}$ of the Ising model and the percolation critical point $p_{c}$ for the FK model with $q=2$ are given by

$$
\begin{equation*}
1-e^{-\beta_{c}}=p_{c}=\frac{\sqrt{2}}{1+\sqrt{2}} . \tag{2.4}
\end{equation*}
$$

The dual lattice $\left(\mathbb{Z}^{2}\right)^{*}$ is $\mathbb{Z}^{2}$ shifted by $\left(\frac{1}{2}, \frac{1}{2}\right)$; sites and bonds of this lattice are called dual sites and dual bonds. $x^{*}$ denotes $x+\left(\frac{1}{2}, \frac{1}{2}\right)$. When necessary for clarity, bonds of $\mathbb{Z}^{2}$ are called regular bonds. To each regular bond $e$ there is associated a unique dual bond $e^{*}$ which is its perpendicular bisector. For $\mathscr{D} \subset \mathscr{B}\left(\mathbb{Z}^{2}\right)$ we write $\mathscr{D}^{*}$ for $\left\{e^{*}: e \in \mathscr{D}\right\}$. For $\Lambda \subset\left(\mathbb{Z}^{2}\right)^{*}, \partial \Lambda$ is defined as for $\Lambda \subset \mathbb{Z}^{2}$, but using adjacency in the dual lattice. The dual bond $e^{*}$ is defined to be open precisely when $e$ is closed, so that for each bond configuration $\omega$ on $\mathbb{Z}^{2}$, there is unique dual configuration $\omega^{*}$ on $\left(\mathbb{Z}^{2}\right)^{*}$. For $p \in[0,1]$ the value $p^{*}$ dual to $p$ at level $q$ is given by

$$
\frac{p}{q(1-p)}=\frac{1-p^{*}}{p^{*}} .
$$

If the regular bonds are distributed as the infinite-volume FK model at $(p, q)$ on $\mathbb{Z}^{2}$ with wired boundary condition, then the dual bonds form the infinite-volume FK model at ( $p^{*}, q$ ) on $\left(\mathbb{Z}^{2}\right)^{*}$ with free boundary condition (see ref. 11.)

An Ising configuration $\sigma \in \Sigma_{\Lambda}$ determines a set of contours, each consisting of dual bonds $e^{*} \in \overline{\mathscr{B}}(\Lambda)^{*}$ for which the corresponding regular bond $e=\langle x y\rangle$ has $\sigma_{x} \neq \sigma_{y}$. In the joint Ising/FK configuration, therefore, contours consist entirely of open dual bonds.

Given sets $\Phi \subset \Lambda$ and a site configuration $\sigma \in \Sigma_{\Lambda}$, we write $\sigma_{\Phi}$ for $\left\{\sigma_{x}: x \in \Phi\right\}$ and let $\mathscr{F}_{\Phi}$ denote the $\sigma$-algebra generated by $\sigma_{\Phi}$. Similarly for $\mathscr{D} \subset \overline{\mathscr{B}}(\Lambda)$ and a bond configuration $\omega \in \Omega_{\Lambda}$, we write $\omega_{\mathscr{D}}$ for $\left\{\omega_{e}: e \in \mathscr{D}\right\}$ and let $\mathscr{G}_{\mathscr{T}}$ denote the $\sigma$-algebra generated by $\omega_{\mathscr{D}}$.

An infinite-volume FK model $P^{p, q}$ (or other bond percolation model) is said to have the weak mixing property if there exist $C, \lambda$ such that, given finite sets $\Phi \subset \Lambda$ and any two bond boundary conditions $\rho_{1}$ and $\rho_{2}$ on $\overline{\mathscr{B}}(\Lambda)^{c}$, we have

$$
\operatorname{Var}\left(P_{\Lambda, \rho_{1}}^{p, q}\left(\omega_{\mathscr{B}(\Phi)} \in \cdot\right), P_{\Lambda, p_{2}}^{p, q}\left(\omega_{\mathscr{B}(\Phi)} \in \cdot\right)\right) \leqslant C \sum_{x \in \Phi, y \notin \Lambda} e^{-\lambda|y-x|},
$$

Loosely this says that the maximum influence, on a fixed region, of the boundary condition decays exponentially to 0 as the boundary recedes to infinity. Equivalently, for all events $A \in \mathscr{G}_{\Lambda^{c}}$ and $B \in \mathscr{G}_{\Phi}$,

$$
\begin{equation*}
\left|P^{p, q}(B \mid A)-P^{p, q}(B)\right| \leqslant C \sum_{x \in \Phi, y \notin A} e^{-\lambda|y-x|} . \tag{2.5}
\end{equation*}
$$

In contrast, $P^{p, q}$ is said to have the ratio weak mixing property if there exist $C, \lambda$ such that, given finite sets $\Phi \subset \Lambda$ and any two bond boundary conditions $\rho_{1}$ and $\rho_{2}$ on $\overline{\mathscr{B}}(\Lambda)^{c}$, we have for all events $A \in \mathscr{G}_{\Lambda^{c}}$ and $B \in \mathscr{G}_{\Phi}$,

$$
\begin{equation*}
\left|\frac{P^{p, q}(A \cap B)}{P^{p, q}(A) P^{p, q}(B)}-1\right| \leqslant C \sum_{x \in \Phi, y \notin A} e^{-\lambda|y-x|}, \tag{2.6}
\end{equation*}
$$

whenever the right side of this inequality is at most 1 . Note that (2.6) is much stronger than (2.5) for $A, B$ for which the probabilities on the left side of (2.5) are much smaller than the right side of (2.5). Weak mixing for the Ising model has a variety of useful consequences, particularly in two dimensions; see ref. 17. It was shown in ref. 3 that for the FK model in two dimensions, exponential decay of either the connectivity (in infinite volume, with wired boundary) or the dual connectivity (in infinite volume, with free boundary) implies ratio weak mixing. In particular, for $q=2$ and $p>p_{c}(2)$, exponential decay of dual connectivity follows from the known properties (see ref. 18) of Gibbs uniqueness and exponential decay of correlations for the Ising model at inverse temperature $\beta^{*}<\beta_{c}(2)$ corresponding to $p^{*}<p_{c}(2)$. Thus we have the following.

Lemma 2.1. Suppose $p>p_{c}(2)$. Then the FK model $P^{p, 2}$ on $\mathscr{B}\left(\mathbb{Z}^{2}\right)$ has the ratio weak mixing property.

The following is an immediate consequence of the definition of ratio weak mixing.

Lemma 2.2 (ref. 5). Suppose that the FK model $P^{p, q}$ has the ratio weak mixing property. There exists a constant $c_{1}$ as follows. Suppose $r>3$ and $\mathscr{D}, \mathscr{E} \subset \mathscr{B}\left(\mathbb{Z}^{2}\right)$ with $\operatorname{diam}(\mathscr{E}) \leqslant r$ and $d(\mathscr{D}, \mathscr{E}) \geqslant c_{1} \log r$. Then for all $A \in \mathscr{G}_{\mathscr{D}}$ and $B \in \mathscr{G}_{\mathscr{\delta}}$, we have

$$
\frac{1}{2} P^{p, 2}(A) P^{p, 2}(B) \leqslant P^{p, 2}(A \cap B) \leqslant 2 P^{p, 2}(A) P^{p, 2}(B) .
$$

We write $y \stackrel{*}{\leftrightarrow} z$ for the event that $y$ is connected to $z$ by a path of open dual bonds. For $q \geqslant 1, P^{p, q}$ has the FKG property (see ref. 11), so $-\log P^{p, q}\left(0^{*} \stackrel{*}{\leftrightarrow} x^{*}\right)$ is a subadditive function of $x$, and therefore the limit

$$
\begin{equation*}
\tau(x)=\lim _{n \rightarrow \infty}-\frac{1}{n} \log P^{p, q}\left(0^{*} \stackrel{*}{\leftrightarrow}(n x)^{*}\right), \tag{2.7}
\end{equation*}
$$

exists for $x \in \mathbb{Q}^{2}$, provided we take the limit through values of $n$ for which $n x \in \mathbb{Z}^{2}$. This definition extends to $\mathbb{R}^{2}$ by continuity (see ref. 2 ); the result$\operatorname{ing} \tau$ is a norm on $\mathbb{R}^{2}$, when the dual connectivity decays exponentially (i.e., $\tau(x)$ is positive for all $x \neq 0$, or equivalently by lattice symmetry, $\tau(x)$ is positive for some $x \neq 0$; we abbreviate this by saying $\tau$ is positive.) By standard subadditivity results,

$$
\begin{equation*}
P^{p, q}\left(0^{*} \stackrel{*}{\mapsto} x^{*}\right) \leqslant e^{-\tau(x)} \quad \text { for all } x . \tag{2.8}
\end{equation*}
$$

In the opposite direction, it is $\mathrm{known}^{(5)}$ that if $\tau$ is positive (so ratio weak mixing holds), then for some $\epsilon_{1}$ and $c_{2}$,

$$
\begin{equation*}
P^{p, q}\left(0^{*} \stackrel{*}{\leftrightarrow} x^{*}\right) \geqslant \epsilon_{1}|x|^{-c_{2}} e^{-\tau(x)} \quad \text { for all } \quad x \neq 0 . \tag{2.9}
\end{equation*}
$$

It follows from the fact that the surface tension $\tau$ is a norm on $\mathbb{R}^{2}$ with axis symmetry that, letting $e_{i}$ denote the $i$ th unit coordinate vector, we have

$$
\begin{equation*}
\frac{1}{\sqrt{2}} \tau\left(e_{1}\right) \leqslant \frac{\tau(x)}{|x|} \leqslant \sqrt{2} \tau\left(e_{1}\right) \quad \text { for all } \quad x \neq 0 \tag{2.10}
\end{equation*}
$$

A weakness of Lemma 2.2 is that the locations $\mathscr{D}, \mathscr{E}$ of the two events must be deterministic. The next lemma from ref. 4 applies only to a limited class of events but allows the locations to be partially random. For $\mathscr{C} \subset \mathscr{D} \subset \mathscr{B}\left(\mathbb{Z}^{2}\right)$ we say an event $A \subset\{0,1\}^{\mathscr{D}}$ occurs on $\mathscr{C}$ (or on $\mathscr{C}^{*}$ ) in $\omega \in\{0,1\}^{\mathscr{D}}$ if $\omega^{\prime} \in A$ for every $\omega^{\prime} \in\{0,1\}^{\mathscr{D}}$ satisfying $\omega_{e}^{\prime}=\omega_{e}$ for all $e \in \mathscr{C}$. For a possibly random set $\mathscr{F}(\omega)$ we say $A$ occurs only on $\mathscr{F}$ (or equivalently, on $\mathscr{F}^{*}$ ) if $\omega \in A$ implies $A$ occurs on $\mathscr{F}(\omega)$ in $\omega$. We say events $A$ and $B$ occur at separation $r$ in $\omega$ if there exist $\mathscr{C}, \mathscr{E} \subset \mathscr{D}$ with $d(\mathscr{C}, \mathscr{E}) \geqslant r$ such that $A$ occurs on $\mathscr{C}$ and $B$ occurs on $\mathscr{E}$ in $\omega$. Let $A \circ_{r} B$ denote the event that $A$ and $B$ occur at separation $r$. Let $\mathscr{D}^{r}=\left\{e \in \mathscr{B}\left(\mathbb{Z}^{2}\right): d(e, \mathscr{D}) \leqslant r\right\}$.

For $x$ a (regular or dual) site, we write $C_{x}=C_{x}(\omega)$ for the (regular or dual) cluster of $x$ in the bond configuration $\omega$.

Lemma 2.3 (ref. 4). Let $P^{p, q}$ be an FK model on $\mathscr{B}\left(\mathbb{Z}^{2}\right)$, with $\tau$ positive and $q \geqslant 1$, satisfying the ratio weak mixing property. There exist
constants $c_{i}, \epsilon_{i}$ as follows. Let $\mathscr{D} \subset \mathscr{B}\left(\mathbb{Z}^{2}\right), x \in\left(\mathbb{Z}^{2}\right)^{*}$ and $r>c_{3} \log |\mathscr{D}|$, and let $A, B$ be events such that $A$ occurs only on $C_{x}$ and $B \in \mathscr{G}_{\mathscr{Q}}$. Then

$$
\begin{equation*}
P^{p, q}\left(A \circ_{r} B\right) \leqslant\left(1+c_{4} e^{-\epsilon_{2} r}\right) P^{p, q}(A) P^{p, q}(B) . \tag{2.11}
\end{equation*}
$$

Let $\overline{x y}$ denote the line through $x$ and $y$. Let

$$
H_{a}^{+}=\left\{\left(x_{1}, x_{2}\right) \in \mathbb{R}^{2}: x_{2} \geqslant a\right\}
$$

and let $H_{x y}$ denote the closed halfspace bounded by $\overline{x y}$ which is to the right as one moves from $x$ to $y$.

Lemma 2.4 (ref. 5). Let $P^{p, q}$ be an FK model on $\mathscr{B}\left(\mathbb{Z}^{2}\right)$, with $\tau$ positive and $q \geqslant 1$, satisfying the ratio weak mixing property. There exist $\epsilon_{3}, c_{5}$ such that for all $x \neq y \in \mathbb{R}^{2}$ and all dual sites $u, v \in H_{x y}$,

$$
P\left(u \stackrel{*}{\leftrightarrow} v \text { via a path in } H_{x y}\right) \geqslant \frac{\epsilon_{3}}{|v-u|^{c_{5}}} e^{-\tau(v-u)} .
$$

For the $\tau$ of (2.7), with $q=2$, the next lemma is an easy consequence of the sharp triangle inequality satisfied by $\tau$ (see ref. 14), which is obtained from the exact solution of the Ising model on $\mathbb{Z}^{2}$. We will not use this method here, however, to help make it apparent that our results are not specific to the Ising model.

For $u \in \mathbb{R}^{2}$ let $D_{u}^{+}$and $D_{u}^{-}$denote the diagonal lines through $u$ and $u+(1,1)$ and through $u$ and $u+(1,-1)$, respectively. Let $T_{N, k}^{1}$ denote the triangle with vertices $(-k,-N),(k,-N)$ and $(0,-(N-k)$ ). Note the base of $T_{N, k}^{1}$ is in the bottom side of $\Lambda_{N}$ and the other two sides are parallel to the diagonals. Let $T_{N, k}^{i}, i=2,3,4$, be the corresponding triangles (obtained by rotation) with bases in the left, top and right sides of $\Lambda_{N}$, respectively.

Lemma 2.5. Suppose $\tau$ is a norm on the plane which has axis and diagonal symmetry. There exists a constant $\epsilon_{4}$ as follows. Let $0<k$ $<k+2 m<N$ and $x=\left(x_{1},-N-\frac{1}{2}\right), y=\left(y_{1},-N-\frac{1}{2}\right)$ with $x_{1} \leqslant-k, y_{1} \geqslant k$, and $z=\left(z_{1}, z_{2}\right) \in H_{-N-\frac{1}{2}}^{+} \backslash T_{N+\frac{1}{2}}^{1}, k+2 m+\frac{1}{2}$. Then

$$
\begin{equation*}
\tau(z-x)+\tau(y-z) \geqslant 2 k \tau\left(e_{1}\right)+\epsilon_{4} m . \tag{2.12}
\end{equation*}
$$

Proof. If $x_{1} \leqslant-k-m$ or $y_{1} \geqslant k+m$ then $\tau(y-x) \geqslant \tau\left((2 k+m) e_{1}\right)$ and (2.12) follows easily. Hence assume $\left|x_{1}\right|, y_{1} \in[k, k+m)$. We may also assume $z_{1} \geqslant 0$.

If $z_{1}>y_{1}$ then from symmetry and convexity we have $\tau(z-x) \geqslant$ $\tau(y-x)$ and $\tau(y-z) \geqslant \tau\left(m e_{1} / 2\right)$, and (2.12) follows easily. Hence we assume $0 \leqslant z_{1} \leqslant y_{1}$.

If $z$ is above $D_{x}^{+}$, let $z^{\prime}$ be the reflection of $z$ across $D_{x}^{+}$. Then $z^{\prime} \notin T_{N, k+2 m}^{1}$ and $\tau\left(z^{\prime}-x\right)=\tau(z-x)$, and it follows easily from symmetry and convexity that $\tau(y-z) \geqslant \tau\left(y-z^{\prime}\right)$. Thus it is sufficient to prove (2.12) for $z^{\prime}$. Hence we may assume $z$ is on or below $D_{x}^{+}$.

Now let $u=\left(u_{1}, u_{2}\right)$ be the reflection of $z$ across $D_{y}^{-}$, and let $v=\left(v_{1}, v_{2}\right)$ be the point where $\overline{x z}$ intersects $D_{y}^{-}$. By the above assumptions on $z$ and simple geometry, we have $x_{1} \leqslant u_{1} \leqslant v_{1}$ and $x_{2} \leqslant u_{2} \leqslant v_{2}$. Using symmetry and convexity we therefore obtain

$$
\tau(z-x)=\tau(z-v)+\tau(v-x) \geqslant \tau(z-v)+\tau(u-x) .
$$

Since $\tau(y-z)=\tau(y-u)$, it follows that

$$
\begin{aligned}
\tau(z-x)+\tau(y-z) & \geqslant \tau(z-v)+\tau(u-x)+\tau(y-u) \\
& \geqslant \epsilon_{4} m+\tau(y-x),
\end{aligned}
$$

as desired.
For $x, y \in\left(\mathbb{Z}^{2}\right)^{*}, r>0$ and $G \subset \mathbb{R}^{2}$, we say there is an $r$-near dual connection from $x$ to $y$ in $G$ if for some $u, v \in\left(\mathbb{Z}^{2}\right)^{*}$ with $d(u, v) \leqslant r$, there are open dual paths from $x$ to $u$ and from $y$ to $v$ in $G$. Let $N(x, y, r, G)$ denote the event that such an $r$-near dual connection exists. The following result is from ref. 5.

Lemma 2.6. Let $P^{p, q}$ be an FK model on $\mathscr{B}\left(\mathbb{Z}^{2}\right)$, with $q \geqslant 1$, for which $\tau$ is positive. There exist $c_{i}$ such that if $|x|>1$ and $r \geqslant c_{6} \log |x|$ then

$$
P^{p, q}\left(N\left(0, x, r, \mathbb{R}^{2}\right)\right) \leqslant e^{-\tau(x)+c \gamma^{r}} .
$$

The next lemma shows that a dual connection via a site for which the triangle inequality is strict by an amount $t>0$ has an excess cost proportional to $t$.

Lemma 2.7. Let $P^{p, q}$ be an FK model on $\mathscr{B}\left(\mathbb{Z}^{2}\right)$, with $q \geqslant 1$, for which $\tau$ is positive. There exists $c_{8}$ as follows. Suppose $x, y, z \in\left(\mathbb{Z}^{2}\right)^{*}$ and $t \geqslant c_{8} \log |y-x|$ satisfy $|y-x|>1$ and

$$
\tau(z-x)+\tau(y-z) \geqslant \tau(y-x)+t .
$$

Then

$$
P^{p, q}(x \stackrel{*}{\longleftrightarrow} z \stackrel{*}{\hookrightarrow} y) \leqslant e^{-\tau(y-x)-\frac{1}{20} t} .
$$

Proof. By Lemma 2.1, $P^{p, q}$ has the ratio weak mixing property. Let $B=B_{\tau}(z, 3(\tau(y-x)+t))$. Then provided $c_{8}$ is large,

$$
\begin{equation*}
P^{p, q}\left(z \stackrel{*}{\longleftrightarrow} B^{c}\right) \leqslant c_{9}(\tau(y-x)+t) e^{-\tau(y-x)-t} \leqslant \frac{1}{3} e^{-\tau(y-x)-\frac{1}{2} t} \tag{2.13}
\end{equation*}
$$

Thus we need only consider paths inside $B$. Let $\epsilon_{5}>0$ to be specified and consider $\omega \in[x \stackrel{*}{\leftrightarrow} z \stackrel{*}{\leftrightarrow} y] \cap N\left(x, y, \epsilon_{5} t, B \backslash B_{\tau}(z, t / 5)\right)^{c}$. For such $\omega$, there exist $z^{\prime}, z^{\prime \prime} \in \partial\left(B_{\tau}(z, t / 5) \cap\left(\mathbb{Z}^{2}\right)^{*}\right)$ and paths $x \stackrel{*}{\leftrightarrow} z^{\prime}, y \stackrel{*}{\leftrightarrow} z^{\prime \prime}$ in $B$ occuring at separation $\epsilon_{5} t$. Now

$$
\tau\left(z^{\prime}-x\right)+\tau\left(y-z^{\prime \prime}\right) \geqslant \tau(x-x)+\tau(y-z)-2\left(\frac{t}{5}+\tau\left(e_{1}\right)\right) \geqslant \tau(y-x)+\frac{1}{2} t,
$$

so using Lemma 2.3, provided $c_{8}$ is large,

$$
\begin{align*}
P^{p, q} & \left([x \stackrel{*}{\leftrightarrow} z \stackrel{*}{\leftrightarrow} y \text { in } B] \cap N\left(x, y, \epsilon_{5} t, B \backslash B_{\tau}(z, t / 5)\right)^{c}\right) \\
& \leqslant \sum_{z^{\prime}, z^{\prime \prime}} 2 P^{p, q}\left(x \stackrel{*}{\leftrightarrow} z^{\prime}\right) P^{p, q}\left(z^{\prime \prime} \leftrightarrow y\right) \\
& \leqslant 2\left|\partial\left(B_{\tau}(z, t / 5) \cap\left(\mathbb{Z}^{2}\right)^{*}\right)\right|^{2} e^{-\tau(y-x)-t / 2} \\
& \leqslant \frac{1}{3} e^{-\tau(y-x)-t / 4} . \tag{2.14}
\end{align*}
$$

Next, provided $\epsilon_{5}$ is small, an application of Lemma 2.2 gives

$$
\begin{align*}
P^{p, q} & \left([x \stackrel{*}{\leftrightarrow} z \stackrel{*}{\leftrightarrow} y \text { in } B] \cap N\left(x, y, \epsilon_{5} t, B \backslash B_{\tau}(z, t / 5)\right)\right) \\
& \leqslant P^{p, q}\left(\left[z \stackrel{*}{\hookrightarrow} B_{\tau}(z, t / 10)^{c}\right] \cap N\left(x, y, \epsilon_{5} t, B \backslash B_{\tau}(z, t / 5)\right)\right) \\
& \leqslant 2 P^{p, q}\left(z \stackrel{*}{\hookrightarrow} B_{\tau}(z, t / 10)^{c}\right) P^{p, q}\left(N\left(x, y, \epsilon_{5} t, B \backslash B_{\tau}(z, t / 5)\right)\right) \\
& \leqslant 2 c_{10} t e^{-t / 10} e^{-\tau(y-x)+c_{7} \epsilon_{5} t} \\
& \leqslant \frac{1}{3} e^{-\tau(y-x)-t / 20} \tag{2.15}
\end{align*}
$$

Together, (2.13), (2.14) and (2.15) complete the proof.

## 3. PROOF OF THEOREM 1.1(i)

Let

$$
\bar{\Lambda}=\Lambda \cup \partial \Lambda, \quad \Lambda \subset \mathbb{Z}^{2}
$$

By a plus path in a site configuration $\sigma$ we mean a lattice path on which all sites $z$ have $\sigma_{z}=1$; minus paths are defined analogously. We write $x \stackrel{+}{\leftrightarrows} y$ $(x \leftrightarrow y)$ for the event that $x$ is connected to $y$ by a plus (minus) path. For $\Phi \subset \bar{\Lambda}_{N}$, the cluster of $\Phi$ in a bond configuration $\omega \in\{0,1\}^{\overline{\mathscr{G}}\left(\Lambda_{N}\right)}$ is the set

$$
C(\Phi, \omega)=\left\{x \in \Lambda: \text { in } \omega, x \leftrightarrow \Phi \text { in } \overline{\mathscr{B}}\left(\Lambda_{N}\right)\right\} .
$$

The plus cluster of $\Phi$ in a site configuration $\sigma \in \Sigma_{A_{N}}$ is the set

$$
C_{+}(\Phi, \sigma)=\left\{x \in \Lambda: \text { in } \sigma, x \stackrel{+}{\leftrightarrows} \Phi \text { in } \overline{\mathscr{B}}\left(\Lambda_{N}\right)\right\} .
$$

If $\sigma_{x}=-1$, then of course $C_{+}(x, \sigma)$ is empty. The minus cluster $C_{-}(\Phi, \sigma)$ is defined analogously.

For $\Phi \subset \mathbb{Z}^{2}$ we define

$$
Q(\Phi)=\bigcup_{x \in \Phi}\left(x+\left[-\frac{1}{2}, \frac{1}{2}\right]^{2}\right) .
$$

Here $x+\left[-\frac{1}{2}, \frac{1}{2}\right]^{2}$ denotes the translation of $\left[-\frac{1}{2}, \frac{1}{2}\right]^{2}$ by $x$. For $x \in \mathbb{R}^{2}$ and $r>0$ we let $B(x, r)$ and $B_{\tau}(x, r)$ be the closed Euclidean ball and $\tau$-ball, respectively, of radius $r$ about $x$.

Fix $\beta>\beta_{c}, N \geqslant 1$ and $K \log N \leqslant k \leqslant N$, with $K$ to be specified later. Let $p=1-e^{-\beta}$ and $j=(N-k) / 4$. We assume $k$ and $j$ are integers; the modifications otherwise are trivial. Let $\mathbb{P}_{N, k}^{\epsilon}$ denote the joint site/bond distribution obtained using the percolation construction of the FK model, for which the site marginal distribution is $\mu_{\Lambda_{N}, \eta^{k, \epsilon}}^{\beta}$ and the bond marginal distribution is $P_{A_{N}, \eta^{k, \epsilon}}^{p, 2}$. (To avoid ambiguous notation we write $\mathbb{P}_{N, k}^{-}$for $\mathbb{P}_{N, k \cdot}^{-1}$.) We write $(\sigma, \omega)$ for a generic joint configuration in $\Sigma_{\Lambda_{N}} \times\{0,1\}^{\overline{\mathcal{F}}\left(\Lambda_{N}\right)}$. We call $(\sigma, \omega)$ allowable (under $\left.\eta^{k, \epsilon}\right)$ if $\mathbb{P}_{N, k}^{\epsilon}((\sigma, \omega))>0$. Define the strip of sites $\Gamma_{1}$ by $\Gamma_{1}=([-k, k] \times\{-N-1\}) \cap \mathbb{Z}^{2}$, and let $\Gamma_{i}, i=2,3,4$, be the corresponding strips of sites, obtained by rotation, in the left, top and right sides of $\partial \tilde{\Lambda}_{N+1}$, respectively. (We will refer to the side corresponding to subscript $i$ as the $i$ th side of $\tilde{\Lambda}_{n}$, for general $n$.) For $i=1,2,3$ let $\Gamma_{i, i+1}$ be the set of sites in $\partial \Lambda_{N}$ which are between $\Gamma_{i}$ and $\Gamma_{i+1}$, in the obvious sense, and let $\Gamma_{4,5}=\Gamma_{4,1}=\Gamma_{0,1}$ be the set of sites in $\partial \Lambda_{N}$ which are between $\Gamma_{4}$ and $\Gamma_{1}$. We also include the appropriate "corner site" as an element of $\Gamma_{i, i+1}$, e.g., $(-N-1,-N-1) \in \Gamma_{1,2}$. Let

$$
\Gamma_{i, m}^{*}=\left\{x \in\left(\mathbb{Z}^{2}\right)^{*} \cap \partial \tilde{\Lambda}_{N+\frac{1}{2}}: x \text { is a corner of } Q(y) \text { for some } y \in \Gamma_{i, m}\right\} .
$$

Let $D_{i}$ be the event that there is no plus-path in $\sigma$ in $\overline{\mathscr{B}}\left(\Lambda_{N}\right)$ from $\Gamma_{i}$ to $\left(T_{N, k+3 j}^{i}\right)^{c}$, and $D=\bigcap_{i=1}^{4} D_{i}$. In the FK model, an event closely related to $D_{i}$ is

$$
E_{i}=\left\{\omega \in\{0,1\}^{\bar{G}\left(\Lambda_{N}\right)}: \Gamma_{i-1, i}^{*} \stackrel{*}{\mapsto} \Gamma_{i, i+1}^{*} \text { in } T_{N+\frac{1}{2}, k+j+\frac{1}{2}}^{i}\right\} .
$$

We begin with a lower bound on the probability of $E_{i}$. The main point is that restricting the path to lie in $T_{N+\frac{1}{2}, k+j+\frac{1}{2}}^{i}$ does not excessively alter the probability of an open dual path from $\Gamma_{i-1, i}^{*}$ to $\Gamma_{i, i+1}^{*}$.

Lemma 3.1. Let $p>p_{c}(2)$. There exist $c_{i}, \epsilon_{i}$ such that for $N, k \geqslant 1$ with $c_{11} \log N \leqslant N-k \leqslant N$, and $E_{i}$ as above,

$$
P_{\Lambda_{N}, \eta^{k, 0}}^{p, 2}\left(E_{i}\right) \geqslant \frac{\epsilon_{6}}{k^{c_{13}}} e^{-2 k \tau\left(e_{1}\right)}
$$

Proof. We may assume $i=1$ and $k \geqslant c_{12}$, with $c_{12}$ to be specified. Let $m=c_{14} \log k$ and $n=\left\lfloor c_{15} \log k\right\rfloor$, where $c_{14}>c_{15}$ are to be specified and $\lfloor\cdot\rfloor$ denotes the integer part. Provided $c_{11}$ is large (depending on $c_{14}$ ), we have $m \leqslant j$. Let $x^{\prime}=\left(-k-\frac{1}{2},-N-\frac{1}{2}\right), y^{\prime}=\left(k+\frac{1}{2},-N-\frac{1}{2}\right), x=x^{\prime}+(n, n), y=$ $y^{\prime}+(-n, n)$ and let

$$
E_{1}^{\prime}=\left\{\omega \in\{0,1\}^{\bar{g}(1)}: x \stackrel{*}{\leftrightarrow} y \text { in } T_{N, k+2 m}^{1} \cap H_{-N+n}^{+}\right\} .
$$

By Lemma 2.4, for some $\epsilon_{7}$,

$$
\begin{equation*}
P^{p, 2}\left(x \stackrel{*}{\mapsto} y \text { in } H_{-N+n}^{+}\right) \geqslant \frac{\epsilon_{7}}{k^{c_{5}}} e^{-2 k t\left(e_{1}\right)} . \tag{3.1}
\end{equation*}
$$

Note this probability is for the infinite-volume limit. For each dual site $z \in H_{-N+n}^{+} \backslash T_{N, k+2 m}^{1}$, let $F_{z}$ denote the event that there exist open dual paths $x^{\prime} \stackrel{*}{\leftrightarrow} z \stackrel{*}{\leftrightarrow} y^{\prime}$. By (3.1),

$$
\begin{equation*}
P^{p, 2}\left(E_{1}^{\prime}\right) \geqslant \frac{\epsilon_{7}}{k^{c_{5}}} e^{-2 k t\left(e_{1}\right)}-P^{p, 2}\left(\bigcup_{z \in\left(\mathbb{Z}^{2}\right)^{*} \cap H_{-N+n}^{+} \backslash T_{N, k+2 m}^{1}} F_{z}\right) . \tag{3.2}
\end{equation*}
$$

We wish to show that the second term on the right side of (3.2) is at most half of the first term on the right side. Let $\Theta=\left(\mathbb{Z}^{2}\right)^{*} \cap B_{\tau}\left(x, 3 k \tau\left(e_{1}\right)\right) \cap$ $H_{-N+n}^{+} \backslash T_{N, k+2 m}^{1}$. We have

$$
\begin{equation*}
P^{p, 2}\left(\bigcup_{z \in\left(\mathbb{Z}^{2}\right)^{*} \cap H_{-N+n}^{+} \backslash T_{N, k+2 m}^{1}} F_{z}\right) \leqslant P^{p, 2}\left(x \stackrel{*}{\leftrightarrow} B_{\tau}\left(x, 3 k \tau\left(e_{1}\right)\right)^{c}\right)+\sum_{z \in \Theta} P^{p, 2}\left(F_{z}\right) . \tag{3.3}
\end{equation*}
$$

By (2.8), provided $c_{12}$ is large,

$$
\begin{equation*}
P^{p, 2}\left(x \stackrel{*}{\leftrightarrows} B_{\tau}\left(x, 3 k \tau\left(e_{1}\right)\right)^{c}\right) \leqslant c_{16} k e^{-3 k \tau\left(e_{1}\right)} \leqslant \frac{1}{4} \frac{\epsilon_{7}}{k^{c_{5}}} e^{-2 k \tau\left(e_{1}\right)} . \tag{3.4}
\end{equation*}
$$

Fix $z \in \Theta$. We decompose the event $F_{z}$ according to whether there is a $\left(c_{6} \log 5 k\right)$-near dual connection from $x$ to $y$ in $B\left(z, 3 c_{17} \log k\right)^{c}$, where $c_{6}$ is from Lemma 2.6 and $c_{17}$ is to be specified. Let $\Psi_{z}=B\left(z, c_{17} \log k\right) \cap\left(\mathbb{Z}^{2}\right)^{*}$. Using (2.8) and Lemma 2.2, provided $c_{17}$ and $c_{12}$ are large we obtain

$$
\begin{align*}
P^{p, 2}( & \left(F_{z} \cap N\left(x, y, c_{6} \log 5 k, B\left(z, 2 c_{17} \log k\right)^{c}\right)\right) \\
& \leqslant P^{p, 2}\left(\left[z \stackrel{*}{\leftrightarrows} \partial \Psi_{z}\right] \cap N\left(x, y, c_{6} \log 5 k, B\left(z, 2 c_{17} \log k\right)^{c}\right)\right) \\
& \leqslant 2 P^{p, 2}\left(z \stackrel{*}{\leftrightarrow} \partial \Psi_{z}\right) P^{p, 2}\left(N\left(x, y, c_{6} \log 5 k, B\left(z, 2 c_{17} \log k\right)^{c}\right)\right) \\
& \leqslant 2 c_{18} k e^{-\frac{1}{2} c_{17} \tau\left(e_{1}\right) \log k} e^{-\tau(y-x)+c c_{6} \log 5 k} \\
& \leqslant e^{-\frac{1}{4} c_{17} \tau\left(e_{1}\right) \log k-2(k-n) \tau\left(e_{1}\right)}, \tag{3.5}
\end{align*}
$$

where $c_{7}$ is from Lemma 2.6. Since $|\Theta| \leqslant c_{19} k^{2}$, provided we choose $c_{17}$ large enough (depending on $c_{15}$ ) this gives

$$
\begin{equation*}
\sum_{z \in \Theta} P^{p, 2}\left(F_{z} \cap N\left(x, y, c_{6} \log 5 k, B\left(z, 2 c_{17} \log k\right)^{c}\right)\right) \leqslant \frac{1}{8} \frac{\epsilon_{7}}{k^{c_{5}}} e^{-2 k \tau\left(e_{1}\right)} . \tag{3.6}
\end{equation*}
$$

Next, let $r=c_{6} \log 5 k$ and for $z \in \Theta$ let $\Psi_{z}^{\prime}=B\left(z, 2 c_{17} \log k\right) \cap\left(\mathbb{Z}^{2}\right)^{*}$. We have

$$
\begin{equation*}
F_{z} \cap N\left(x, y, c_{6} \log 5 k, B\left(z, 2 c_{17} \log k\right)^{c}\right)^{c} \subset \bigcup_{u, v \in \partial \Psi_{z}^{\prime}}\left([x \leftrightarrow u] \circ_{r}[y \leftrightarrow v]\right) . \tag{3.7}
\end{equation*}
$$

Now for $u, v \in \partial \Psi_{z}^{\prime}$,

$$
\tau(u-x) \geqslant \tau(z-x)-\tau(u-z) \geqslant \tau(z-x)-3 c_{17} \tau\left(e_{1}\right) \log k
$$

and similarly

$$
\tau(y-v) \geqslant \tau(y-z)-3 c_{17} \tau\left(e_{1}\right) \log k .
$$

Hence provided $c_{14}$ is large enough (depending on $c_{17}$ ), we obtain using Lemma 2.5 that

$$
\begin{aligned}
\tau(u-x)+\tau(y-v) & \geqslant 2(k-n) \tau\left(e_{1}\right)+\epsilon_{4} m-6 c_{17} \tau\left(e_{1}\right) \log k \\
& \geqslant 2(k-n) \tau\left(e_{1}\right)+\frac{\epsilon_{4}}{2} m .
\end{aligned}
$$

Combining this with (3.7), Lemma 2.3 and (2.8), provided $c_{14}$ and $c_{14} / c_{15}$ are large we get

$$
\begin{align*}
P^{p, 2} & \left(F_{z} \cap N\left(x, y, c_{6} \log 5 k, B\left(z, 2 c_{17} \log k\right)^{c}\right)^{c}\right) \\
& \leqslant \sum_{u, v \in \partial \Psi_{z}^{\prime}} 2 P^{p, 2}(x \stackrel{*}{\leftrightarrow} u) P^{p, 2}(y \stackrel{*}{\leftrightarrow} v) \\
& \leqslant\left|\partial \Psi_{z}^{\prime}\right|^{2} e^{-2(k-n) \tau\left(e_{1}\right)-\epsilon_{4} m / 2} \\
& \leqslant e^{-2 k \tau\left(e_{1}\right)-\epsilon_{4} m / 4} \tag{3.8}
\end{align*}
$$

and then

$$
\begin{align*}
& \sum_{z \in \Theta} P^{p, 2}\left(F_{z} \cap N\left(x, y, c_{6} \log 5 k, B\left(z, 2 c_{17} \log k\right)^{c}\right)^{c}\right) \\
& \quad \leqslant c_{19} k^{2} e^{-2 k \tau\left(e_{1}\right)-\epsilon_{4} m / 4} \leqslant \frac{1}{8} \frac{\epsilon_{7}}{k^{c}} e^{-2 k \tau\left(e_{1}\right)} . \tag{3.9}
\end{align*}
$$

Combining (3.2), (3.3), (3.6) and (3.9) we obtain

$$
P^{p, 2}\left(E_{1}^{\prime}\right) \geqslant \frac{1}{2} \frac{\epsilon_{7}}{k^{c_{5}}} e^{-2 k \tau\left(e_{1}\right)}
$$

Then from Lemma 2.2, provided $c_{15}$ is large,

$$
P_{\Lambda_{N}, \eta^{k, 0}}^{p, 2}\left(E_{1}^{\prime}\right) \geqslant \frac{1}{4} \frac{\epsilon_{7}}{k^{c_{5}}} e^{-2 k \tau\left(e_{1}\right)} .
$$

Let $\gamma_{x}$ and $\gamma_{y}$ be dual paths of (minimal) length $2 n$ from $x$ to $x^{\prime}$ and from $y$ to $y^{\prime}$, respectively, in $T_{N, k+2 m}^{1}$. Let $E_{1}^{\prime \prime}$ denote the event that all dual bonds in $\gamma_{x}$ and $\gamma_{y}$ are open. From the FKG inequality,

$$
\begin{aligned}
& P_{A_{N}, \eta^{k}}^{p, 2},\left(E_{1}\right) \geqslant P_{A_{N}, \eta^{k, 0}}^{p, 2}\left(E_{1}^{\prime} \cap E_{1}^{\prime \prime}\right) \\
& \geqslant P_{A_{N}, \eta^{k}, 0}^{p, 2}\left(E_{1}^{\prime}\right) P_{\Lambda_{N}, \eta^{k}, 0}^{p, 2}\left(E_{1}^{\prime \prime}\right) \\
& \geqslant \frac{1}{4} \frac{\epsilon_{7}}{k^{c_{5}}} e^{-2 k t\left(e_{1}\right)} \alpha(p, 2)^{4 n}
\end{aligned}
$$

and the lemma follows.
For $\omega \in E_{i}, \partial Q\left(C\left(\Gamma_{i}, \omega\right)\right)$ includes a unique open dual path $\gamma_{i}(\omega)$ in $T_{N+\frac{1}{2}, k+j+\frac{1}{2}}^{i}$ from $\Gamma_{i-1, i}^{*}$ to $\Gamma_{i, i+1}^{*}$. This path is "closer to $\Gamma_{i}$ " than any other
open dual path in $\tilde{\Lambda}_{N+\frac{1}{2}}$ from $\Gamma_{i-1, i}^{*}$ to $\Gamma_{i, i+1}^{*}$. Further, for fixed $v$ the event [ $\gamma_{i}=v$ ] depends only on the bond/dual bond configuration in the closed region, which we denote $I(v)$, between $v$ and the side of $\partial \tilde{\Lambda}_{N+\frac{1}{2}}$ to which $v$ is attached.

Let $E=\bigcap_{i=1}^{4} E_{i}$, and suppose $\omega \in E$. Let

$$
R(\omega)=\tilde{\Lambda}_{N+\frac{1}{2}} \bigcup_{i=1}^{4} I\left(\gamma_{i}(\omega)\right)
$$

Note $0 \in R(\omega)$, and (under boundary condition $\eta^{k, 0}$ ) all the dual bonds forming $\partial R(\omega)$ are open in $\omega$. The latter means that for fixed $U$, conditionally on $R(\omega)=U$ the configuration on $\mathscr{B}(U)$ is the FK model with free boundary condition.

Let

$$
Y_{N}=\left.\tilde{\Lambda}_{N+\frac{1}{2}}\right|_{i=1} ^{4} T_{N+\frac{1}{2}, k+j+\frac{1}{2}}^{i}, \quad Y_{N}^{\prime}=\left\{x \in Y_{N}: d\left(x, \partial Y_{N}\right) \geqslant j\right\} .
$$

Let $0<h<j$ to be specified, let $w_{12}=(k+2 j,-N+h)$, let $\lambda_{12}$ be the vertical line from $w_{12}$ down to $\partial \tilde{\Lambda}_{N}$ at $(k+2 j,-N)$ and let $\chi_{12}$ be the vertical line from $w_{12}$ up to the diagonal $D_{0}^{-}$at $(k+2 j,-k-2 j)$. Using axis symmetry we obtain 7 more corresponding points $w_{i j}$ and paths $\chi_{i j}, \lambda_{i j}$, for $i=1,2,3,4$ and $j=1,2$, with $w_{i j}$ at distance $h$ from side $i$ of $\tilde{\Lambda}_{N}$.

We want to show that with high probability, there are no open dual paths starting from $Y_{N}^{\prime}$, or from near $\chi_{i j}$, which reach $\partial Y_{N}$. Let $y_{i j 11}, y_{i j 2}$ be the endpoints of the dual bond which is dual to the $l$ th bond of $\chi_{i j}$. Then

$$
d\left(y_{i j l m}, \partial Y_{N}\right) \geqslant h+\frac{l}{\sqrt{2}} \quad \text { for all } i, j, l, m
$$

For $x \in Y_{N}$ let

$$
G_{x}=B\left(x, \frac{1}{2} d\left(x, \partial Y_{N}\right)\right) \cap\left(\mathbb{Z}^{2}\right)^{*},
$$

and define

$$
\begin{aligned}
\Psi_{N}= & \left(Y_{N}^{\prime} \cap\left(\mathbb{Z}^{2}\right)^{*}\right) \\
& \cup\left\{y_{i j m}: 1 \leqslant i \leqslant 4 ; j=1,2 ; 1 \leqslant l \leqslant N-h-k-2 j ; m=1,2\right\} .
\end{aligned}
$$

Suppose $U \supset Y_{N}$. Then $d\left(G_{x}, \partial U\right) \geqslant j / 2$ for all $x \in Y_{N}^{\prime}$. Hence using Lemma 2.2, provided $K$ and $h$ are large enough we get

$$
\begin{align*}
P_{\Lambda_{N}, \eta^{k}}^{p, 2}(x & \left.\stackrel{*}{\leftrightarrow} \partial G_{x} \text { for some } x \in \Psi_{N} \mid \omega \in E, R(\omega)=U\right) \\
& =P_{U \cap \mathbb{Z}^{2}, f}^{p, 2}\left(x \stackrel{*}{\leftrightarrow} \partial G_{x} \text { for some } x \in \Psi_{N}\right) \\
& \leqslant \sum_{x \in Y_{N}^{\prime} \cap\left(\mathbb{Z}^{2}\right)^{*}} P_{U \cap \mathbb{Z}^{2}, f}^{p, 2}\left(x \stackrel{*}{\leftrightarrow} \partial G_{x}\right)+\sum_{i, j, l, m} P_{U \cap \mathbb{Z}^{2}, f}^{p, 2}\left(y_{i j l m} \stackrel{*}{\leftrightarrow} \partial G_{y_{i j l m}}\right) \\
& \leqslant \sum_{x \in Y_{N}^{\prime} \cap\left(\mathbb{Z}^{2}\right)^{*}} 2 P^{p, 2}\left(x \stackrel{*}{\leftrightarrow} \partial G_{x}\right)+\sum_{i, j, l, m} 2 P^{p, 2}\left(y_{i j l m} \stackrel{*}{\leftrightarrow} \partial G_{y_{i j l m}}\right) \\
& \leqslant\left|Y_{N}^{\prime} \cap\left(\mathbb{Z}^{2}\right)^{*}\right| e^{-j \tau\left(e_{1}\right) / 2}+c_{20} \sum_{l \geqslant 1}(h+l) e^{-\frac{1}{4} \tau\left(e_{1}\right)(h+l)} \\
& \leqslant \frac{1}{2} . \tag{3.10}
\end{align*}
$$

Let $F$ denote the event that $x \stackrel{*}{\leftrightarrow} \partial G_{x}$ for no $x \in \Psi_{N}$, and all bonds in $\lambda_{i j}$ are open for all $i, j$. If $\omega \in E \cap F$, then there is an open circuit in $Y_{N}$ surrounding $\Psi_{N}$ and for each $i$, a portion of this open circuit, together with $\lambda_{i 1}$ and $\lambda_{i 2}$, forms an open path in $T_{N, k+3 j}^{i} \backslash T_{N, k+j}^{i}$ from a site adjacent to $\Gamma_{i-1, i}$ to a site adjacent to $\Gamma_{i, i+1}$. When this occurs (with $\omega \in E \cap F$ ), we call this circuit together with all $\lambda_{i j}$ a blocking pattern. Note the blocking pattern is contained in $R(\omega)$. We have using (3.10) and the FKG inequality that for all $U \supset Y_{N}$,

$$
\begin{align*}
& P_{\Lambda_{N}, \eta^{\prime},,}^{p, 2}(\text { there is a blocking pattern in } R(\omega) \mid \omega \in E, R(\omega)=U) \\
&=P_{U \cap \mathbb{Z}^{2}, f}^{p, 2}(\text { there is a blocking pattern in } U) \\
& \geqslant P_{U \cap \mathbb{Z}^{2}, f}^{p, 2}(F) \\
& \geqslant \frac{1}{2} P_{U \cap \mathbb{Z}^{2}, f}^{p, 2}\left(\text { all bonds in } \lambda_{i j} \text { are open for all } i, j\right) \\
& \geqslant \frac{1}{2} \alpha(p, 2)^{8 h} . \tag{3.11}
\end{align*}
$$

But considering the cluster-labeling construction of the joint Ising/FK configuration, we see that if the configuration $\omega \in E$ has a blocking pattern in $R(\omega)$ and all sites $x$ in the blocking pattern have $\sigma_{x}=-1$ (which occurs with probability $1 / 2$, given such $\omega$ ), then $\sigma \in D$. Thus from (3.11), the FKG property and Lemma 3.1, for some $\epsilon_{8}$,

$$
\begin{equation*}
\mu_{A_{N}, \eta^{k, 0}}^{\beta}(D) \geqslant \frac{1}{4} \alpha(p, 2)^{8 h} P_{A_{N}, \eta^{\prime}, 0}^{p, 2}(E) \geqslant \frac{\epsilon_{8}}{k^{c_{13}}} e^{-8 k \tau\left(e_{1}\right)} . \tag{3.12}
\end{equation*}
$$

We turn now to upper bounds on $\mu_{\Lambda_{N}, \eta^{k}, 0}^{\beta}\left(\partial_{i n} D\right)$. Analogously to $\gamma_{i}(\omega)$, for $\sigma \in D, \partial Q\left(C_{+}\left(\Gamma_{i}, \sigma\right)\right)$ includes a unique open dual path $\gamma_{i}^{+}(\sigma)$ in $T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i}$ from $\Gamma_{i, i+1}^{*}$ to $\Gamma_{i-1, i}^{*}$, for each $i$. For fixed $v$ the event $\left[\gamma_{i}^{+}=v\right.$ ] depends only on the site configuration in $\bar{I}(v) \cap \mathbb{Z}^{2}$.

Suppose $\sigma \in \partial_{i n} D$ and $\sigma^{x} \notin D$. Then for some $i$ we have $x \in \partial C_{+}\left(\Gamma_{i}, \sigma\right)$, and either $x \in \partial T_{N, k+3 j+1}^{i}$ or there is an open dual circuit $\gamma$ in $\omega$ outside $I\left(\gamma_{i}^{+}(\sigma)\right)$ which includes an edge of $Q(x)$ and surrounds some site outside $T_{N, k+3 j}^{i}$. We can choose $\gamma$ to be the outer boundary of $Q(\Phi)$ for some pluscluster $\Phi$ in $\sigma$. In this case we call $\gamma$ an appendable circuit attachable at $x$.

According to Lemma 2.2, we can choose a constant $c_{21}$ as follows. Let $Z_{N}^{i, i+1}=\left\{x \in \mathbb{R}^{2}: d\left(x, \Gamma_{i, i+1}\right) \leqslant c_{21} \log N\right\}$ and $Z_{N}=\bigcup_{i} Z_{N}^{i, i+1}$. Let $V_{\text {closed }}$ be the event that all bonds in $\left\{\langle x y\rangle: y \in \Lambda_{N}, x \in \Gamma_{i, i+1}\right.$ for some $\left.i\right\}$ are closed. Let $V_{\text {open }}$ be the event that all bonds in $\mathscr{B}\left(\mathbb{Z}^{2}\right) \backslash \bar{B}\left(\Lambda_{N}\right)$ are open. (Note the boundary condition $\eta^{k, 0}$ conditions $\omega$ on $V_{\text {open }} \cap V_{\text {closed }}$.) Then for all events $A \in \mathscr{G}_{\bar{G}\left(\Lambda_{N}\right) \backslash \mathscr{B}\left(Z_{N}\right)}$, we have

$$
\begin{equation*}
\frac{1}{2} P^{p, 2}(A) \leqslant P^{p, 2}\left(A \mid V_{\text {closed }}\right) \leqslant 2 P^{p, 2}(A) . \tag{3.13}
\end{equation*}
$$

We therefore call $Z_{N}$ the free-boundary influence region. In particular, provided $K$ is large (depending on $c_{21}$ ), using the FKG inequality and (3.13) we have

$$
\begin{align*}
P_{\Lambda_{N}, \eta^{k, 0}}^{p, 2}\left(Z_{N}^{i-1, i} \stackrel{*}{\leftrightarrow} Z_{N}^{i, i+1}\right) & =P^{p, 2}\left(Z_{N}^{i-1, i} \stackrel{*}{\leftrightarrow} Z_{N}^{i, i+1} \mid V_{\text {open }} \cap V_{\text {closed }}\right) \\
& \leqslant P^{p, 2}\left(Z_{N}^{i-1, i} \stackrel{*}{\leftrightarrow} Z_{N}^{i, i+1} \mid V_{\text {closed }}\right) \\
& \leqslant 2 P^{p, 2}\left(Z_{N}^{i-1, i} \stackrel{*}{\leftrightarrow} Z_{N}^{i, i+1}\right) \\
& \leqslant c_{22}(\log N) e^{-2\left(k-c_{21} \log N\right) \tau\left(e_{1}\right)} \\
& \leqslant N^{c_{23}} e^{-2 k \tau\left(e_{1}\right)} . \tag{3.14}
\end{align*}
$$

Our main task is roughly to show, using Lemma 2.5, that the probability for a connection $\Gamma_{i-1, i}^{*} \stackrel{*}{\leftrightarrows} \Gamma_{i, i+1}^{*}$ (specifically, part of $\gamma_{i}^{+}(\sigma)$ ) which does not stay inside $T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i}$ is smaller than the right side of (3.12) by at least a factor of $e^{-\epsilon_{9} j}$, for some $\epsilon_{9}$. We must decompose the event $\partial_{\text {in }} D$ into several pieces according to the geometry of the sets $C_{+}\left(\Gamma_{i}, \sigma\right)$ and $C\left(\Gamma_{i}, \omega\right)$. The most difficult case is that of leakage along the (free) boundary, in which $\gamma_{i}^{+}(\sigma)$ goes outside $T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i}$ by way of $Z_{N}$.

We define one more special dual path as follows. For $\Phi \subset \Lambda_{N}$ let

$$
\hat{C}(\Phi, \omega)=\left\{x \in \Lambda_{N} \backslash Z_{N}: \text { in } \omega, x \leftrightarrow \Phi \text { in } \overline{\mathscr{B}}\left(\Lambda_{N}\right) \backslash \mathscr{B}\left(Z_{N}\right)\right\} .
$$

If $(\sigma, \omega)$ is allowable and $\sigma \in D$, then $\partial Q\left(\hat{C}\left(\Gamma_{i}, \omega\right)\right)$ includes a unique open dual path from $Z_{N}^{i, i+1}$ to $Z_{N}^{i-1, i}$ in $\overline{\mathscr{B}}\left(\Lambda_{N}\right) \backslash \mathscr{B}\left(Z_{N}\right)$; we denote this path $\hat{\gamma}_{i}(\omega)$. We have

$$
\hat{\gamma}_{i}(\omega) \subset I\left(\gamma_{i}(\omega)\right) \subset I\left(\gamma_{i}^{+}(\omega)\right) \subset T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} .
$$

Let $u_{i}(\omega)$ and $v_{i}(\omega)$ be the starting and ending sites, respectively, of $\hat{\gamma}_{i}(\omega)$ in $Z_{N}^{i, i+1}$ and $Z_{N}^{i-1, i}$, respecitvely. Also define

$$
W_{N}^{i, i+1}=\left\{x \in \mathbb{R}^{2}: d\left(x, \Gamma_{i, i+1}\right) \leqslant 2 \epsilon_{10} j\right\}, \quad W_{N}=\bigcup_{i} W_{N}^{i, i+1},
$$

where $\epsilon_{10}$ is to be specified. Provided $K$ is large enough (depending on $\epsilon_{10}$ and $c_{21}$ ), we have $Z_{N}^{i, i+1} \subset W_{N}^{i, i+1}$.

Case 1. Consider $\sigma \in \partial_{i n} D$, and $\omega$ with $(\sigma, \omega)$ allowable, for which for some $i$ there exist (in order) dual sites $x, z, y \in \hat{\gamma}_{i}(\omega)$ with
$\tau(y-x) \geqslant\left(2 k-2 c_{21} \log N\right) \tau\left(e_{1}\right), \quad \tau(z-x)+\tau(y-z) \geqslant\left(2 k+4 \epsilon_{10} j\right) \tau\left(e_{1}\right)$.

We let $A_{1}$ denote the set of $(\sigma, \omega)$ for which this occurs, let $J_{i}$ denote the set of all $(x, y, z) \in\left(T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} \cap\left(\mathbb{Z}^{2}\right)^{*}\right)^{3}$ for which (3.15) holds, and let $J_{i}^{\prime}$ denote the set of all $(x, y, z) \in J_{i}$ which also satisfy

$$
\tau(y-x) \geqslant\left(2 k+2 \epsilon_{10} j\right) \tau\left(e_{1}\right) .
$$

As in (3.14), provided K is large enough (depending on $\epsilon_{10}$ and $c_{21}$ ), using (3.13), (3.14) and Lemmas 2.2 and 2.7 we get

$$
\begin{align*}
& \mathbb{P}^{0}\left(A_{1}\right) \leqslant \sum_{i=1}^{4} \sum_{(x, y, z) \in J_{i}} P_{\Lambda_{N}, \eta^{k, 0}}^{p, 2}\left(x \stackrel{*}{\longleftrightarrow} z \stackrel{*}{\longleftrightarrow} y \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} \backslash \mathscr{B}\left(Z_{N}\right) ;\right. \\
& \left.Z_{N}^{l-1, l} \stackrel{*}{\longleftrightarrow} Z_{N}^{l, l+1} \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{l} \backslash \mathscr{B}\left(Z_{N}\right) \text { for all } l \neq i\right) \\
& \leqslant \sum_{i=1}^{4} \sum_{(x, y, z) \in J_{i}} 2 P^{p, 2}\left(x \stackrel{*}{\leftrightarrow} z \stackrel{*}{\leftrightarrow} y \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} \backslash \mathscr{B}\left(Z_{N}\right) ;\right. \\
& \left.Z_{N}^{l-1, l} \stackrel{*}{\hookrightarrow} Z_{N}^{l, l+1} \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{l} \backslash \mathscr{B}\left(Z_{N}\right) \text { for all } l \neq i\right) \\
& \leqslant \sum_{i=1}^{4} \sum_{(x, y, z) \in J_{i}} 16 P^{p, 2}\left(x \stackrel{*}{\leftrightarrow} z \stackrel{*}{\leftrightarrow} y \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} \backslash \mathscr{B}\left(Z_{N}\right)\right) \\
& \cdot \prod_{l \neq i} P^{p, 2}\left(Z_{N}^{l-1, l} \stackrel{*}{\longleftrightarrow} Z_{N}^{l, l+1} \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{l} \backslash \mathscr{B}\left(Z_{N}\right)\right) \\
& \leqslant 16\left(\sum_{i=1}^{4} \sum_{(x, y, z) \in J_{i}^{\prime}} e^{-\tau(y-x)}+\sum_{i=1}^{4} \sum_{(x, y, z) \in J_{i} \backslash J_{i}^{\prime}} e^{-\tau(y-x)-\frac{1}{10} \epsilon_{10} j \tau\left(e_{1}\right)}\right) \\
& \times\left(N^{c_{23}} e^{-2 k \tau\left(e_{1}\right)}\right)^{3} \\
& \leqslant 16\left(4\left|J_{1}^{\prime}\right| e^{-\left(2 k+2 \epsilon_{10} j\right) \tau\left(e_{1}\right)}\right. \\
& \left.+4\left|J_{1} \backslash J_{1}^{\prime}\right| e^{-\left(2 k+\frac{1}{10} \epsilon_{10} j-2 c_{21} \log N\right) \tau\left(e_{1}\right)}\right)\left(N^{c_{23}} e^{-2 k \tau\left(e_{1}\right)}\right)^{3} \\
& \leqslant e^{-\left(8 k+\frac{1}{20} \epsilon_{10} j\right) \tau\left(e_{1}\right)} \text {. } \tag{3.16}
\end{align*}
$$

Case 2. Let

$$
R_{N}^{i}=T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} \cap\left(T_{N+\frac{1}{2}, k+j+\frac{1}{2}}^{i} \cup W_{N}\right) .
$$

We may think of $R_{N}^{i}$ as a "triangle with feet." Let $A_{2}$ denote the set of all $(\sigma, \omega) \in \partial_{i n} D \backslash A_{1}$ for which, for some $i$, there exists a dual site $z \in T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} \backslash R_{N}^{i}$ which is either in $\gamma_{i}^{+}(\sigma)$ or in some appendable circuit attachable at some $x \in \partial C_{+}\left(\Gamma_{i}, \sigma\right)$. (In particular, we can choose $z$ with $z \stackrel{*}{\mapsto} \partial\left(B\left(z, \epsilon_{10} j\right) \cap\left(\mathbb{Z}^{2}\right)^{*}\right)$ and $B\left(z, 2 \epsilon_{10} j\right) \subset \tilde{\Lambda}_{N+1}$.) Suppose $(\sigma, \omega) \in A_{2}$. We claim that $B\left(z, 2 \epsilon_{10} j\right) \cap \hat{\gamma}_{i}(\omega)=\phi$. For all $u \in \Gamma_{i, i+1}^{*}$ and $v \in \Gamma_{i-1, i}^{*}$, by Lemma 2.5 we have

$$
\tau(z-u)+\tau(v-z) \geqslant 2 k \tau\left(e_{1}\right)+\epsilon_{4} j .
$$

Therefore for all $u^{\prime} \in Z_{N}^{i, i+1}, v^{\prime} \in Z_{N}^{i-1, i}$ and $z^{\prime} \in B\left(z, 2 \epsilon_{10} j\right)$, provided $K$ is large (depending on $c_{21}$ ) and $\epsilon_{10}$ is small (depending on $\epsilon_{4}$ ),

$$
\begin{aligned}
\tau\left(z^{\prime}-u^{\prime}\right)+\tau\left(v^{\prime}-z^{\prime}\right) & \geqslant 2 k \tau\left(e_{1}\right)+\frac{1}{2} \epsilon_{4} j-2 c_{21}(\log N) \tau\left(e_{1}\right)-4 \epsilon_{10} j \tau\left(e_{1}\right) \\
& \geqslant\left(2 k+4 \epsilon_{10} j\right) \tau\left(e_{1}\right) .
\end{aligned}
$$

Taking $u^{\prime}=u(\omega), v^{\prime}=v(\omega)$ and comparing to (3.15) we see that since $(\sigma, \omega) \notin A_{1}$, we cannot have $z^{\prime} \in \hat{\gamma}_{i}(\omega)$, proving our claim. Therefore as in (3.16),

$$
\begin{align*}
& Z_{N}^{i-1, i} \stackrel{*}{\hookrightarrow} Z_{N}^{i, i+1} \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} \cap B\left(z, 2 \epsilon_{10} j\right)^{c} \backslash \mathscr{B}\left(Z_{N}\right), \\
& Z_{N}^{l-1, l} \stackrel{*}{\hookrightarrow} Z_{N}^{l, l+1} \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{l}\left(\mathscr{B}\left(Z_{N}\right) \text { for all } l \neq i\right) \\
& \leqslant \sum_{i=1}^{4} \sum_{z \in\left(T_{N+\frac{1}{2}}^{i}, k+3 j+\frac{3}{2}, R_{N}^{i}\right) \cap\left(\mathbb{Z}^{2}\right)^{*}} 2 P^{p, 2}\left(z \stackrel{*}{\leftrightarrow} B\left(z, \epsilon_{10} j\right)^{c} \text { in } B\left(z, \epsilon_{10} j+1\right),\right. \\
& Z_{N}^{i-1, i} \stackrel{*}{\leftrightarrow} Z_{N}^{i, i+1} \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} \cap B\left(z, 2 \epsilon_{10} j\right)^{c} \backslash \mathscr{B}\left(Z_{N}\right), \\
& \left.\left.Z_{N}^{l-1, l} \stackrel{*}{\leftrightarrow} Z_{N}^{l, l+1} \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{l}\right) \mathscr{B}\left(Z_{N}\right) \text { for all } l \neq i\right) \\
& \leqslant 32 \sum_{i=1}^{4} \sum_{z \in\left(T_{N+\frac{1}{2}}^{i}, k+3 j+\frac{3}{2} \backslash R_{N}^{i}\right) \cap\left(\mathbb{Z}^{2}\right)^{*}} P^{p, 2}\left(z \stackrel{*}{\leftrightarrow} B\left(z, \epsilon_{10} j\right)^{c}\right) \\
& \times \prod_{l=1}^{4} P^{p, 2}\left(Z_{N}^{l-1, l} \stackrel{*}{\hookrightarrow} Z_{N}^{l, l+1}\right) \\
& \leqslant c_{24}\left|T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{1} \cap\left(\mathbb{Z}^{2}\right)^{*}\right| j e^{-\frac{1}{2} \epsilon_{10} j \tau\left(e_{1}\right)}\left(N^{c_{23}} e^{-2 k t\left(e_{1}\right)}\right)^{4} \\
& \leqslant e^{-\left(8 k+\frac{1}{4} \epsilon_{10} j\right) \tau\left(e_{1}\right)} \text {. } \tag{3.17}
\end{align*}
$$

Case 3. Let $A_{3}=\partial_{\text {in }} D \backslash\left(A_{1} \cup A_{2}\right)$, and suppose $(\sigma, \omega) \in A_{3}$. In this case, plus spins are "leaking along the boundary," in the following sense: for some $i$, we have $\gamma_{i}^{+}(\sigma) \subset R_{N}^{i}$, and either $\gamma_{i}^{+}(\sigma)$ or some appendable circuit $\gamma$ contains a dual site $z \in W_{N}$ at one of the "toes" of $R_{N}^{i}$, that is, in the right or left side of $T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i}$. Note that such a $\gamma$ necessarily has $\gamma \cap T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} \subset R_{N}^{i}$, since $(\sigma, \omega) \notin A_{2}$. We will assume $z$ is in an appendable circuit, attachable at some site $x$; the case of $z \in \gamma_{i}^{+}(\sigma)$ is similar but slightly simpler. From symmetry, we may also assume $i=1$ and $z$ is in the right side of $T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{1}$. We let $A_{3}^{\prime}(x, z)$ denote the event that $A_{3}$ occurs with a specified choice of $x, z$, with $i=1$ and with $z$ in an appendable circuit, and let $A_{3}^{\prime}=\bigcup_{x, z} A_{3}^{\prime}(x, z)$. We define $A_{3}^{\prime \prime}(z)$ and $A_{3}^{\prime \prime}$ analogously for the case of $z \in \gamma_{i}^{+}(\sigma)$.

The Ising model has the following "bounded energy" property:

$$
\begin{equation*}
\mu_{\Lambda_{N}, \eta^{k}, 0}^{\beta}\left(\sigma_{y}=1 \mid \sigma_{w}, w \neq y\right) \geqslant \frac{1}{1+e^{8 \beta}} \quad \text { for all } \quad\left(\sigma_{w}, w \neq y\right) . \tag{3.18}
\end{equation*}
$$

Given a site $x$, let $\omega^{x}$ denote the configuration given by

$$
\omega_{e}^{x}= \begin{cases}0, & \text { if } x \text { is an endpoint of } e \\ \omega_{e}, & \text { otherwise }\end{cases}
$$

Note that if $(\sigma, \omega)$ is allowable, then so is $\left(\sigma^{x}, \omega^{x}\right)$. Let

$$
\begin{aligned}
B_{m}^{\text {trunc }} & =B \cap\left\{\left(y_{1}, y_{2}\right) \in \mathbb{R}^{2}: y_{1}<m\right\} \quad \text { for } \quad m>0, \quad B \subset \mathbb{R}^{2}, \\
\hat{\Gamma}_{1,2} & =([k, k+2 j) \times\{-N\}) \cap \mathbb{Z}^{2}, \\
\psi_{1,2} & =\{k+2 j\} \times\left[-N,-N+2 \epsilon_{10} j\right]
\end{aligned}
$$

and

$$
\hat{C}_{-}\left(\hat{\Gamma}_{1,2}, \sigma^{x}\right)=\left\{y \in\left(\Lambda_{N}\right)_{k+2 j}^{\text {trunc }}: \text { in } \sigma^{x}, y \leftrightarrow \hat{\Gamma}_{1,2} \text { in }\left(\tilde{\Lambda}_{N}\right)_{k+2 j}^{\text {trunc }}\right\} .
$$

For $J \subset \mathbb{Z}^{2}$ define the boundary condition $\eta^{k, 0, J}$ by

$$
\eta_{u}^{k, 0, J}= \begin{cases}1, & \text { if } u \in J  \tag{3.19}\\ \eta_{u}^{k, 0}, & \text { otherwise }\end{cases}
$$

Fix $x, z$ and suppose $(\sigma, \omega) \in A_{3}^{\prime}(x, z)$. There is then a plus path in $\sigma^{x}$ from $\Gamma_{1}$ to $\psi_{1,2}$ in $R_{N}^{1}$, and $\hat{C}_{-}\left(\hat{\Gamma}_{1,2}, \sigma^{x}\right)$ is contained in the region between this plus path and $\partial\left(\Lambda_{N}\right)_{k+2 j}^{\text {trunc }}$. Since

$$
\left|\partial \hat{C}_{-}\left(\hat{\Gamma}_{1,2}, \sigma^{x}\right) \cap \psi_{1,2}\right| \leqslant 2 \epsilon_{10} j,
$$

using the "bounded energy" property (3.18) of the Ising model we have for fixed $J$

$$
\begin{aligned}
& \mathbb{P}_{N, k}^{0}\left((\sigma, \omega) \in A_{3}^{\prime}(x, z), \hat{C}_{-}\left(\hat{\Gamma}_{1,2}, \sigma^{x}\right)=J\right) \\
& \leqslant \mathbb{P}_{N, k}^{0}\left(\hat{C}_{-}\left(\hat{\Gamma}_{1,2}, \sigma^{x}\right)=J \text {, and in } \omega^{x}, \psi_{1,2}-\frac{1}{2} e_{1} \stackrel{*}{\leftrightarrow} Z_{N}^{4,1}\right. \\
& \text { in } \overline{\mathscr{B}}\left(\Lambda_{N} \backslash \bar{J}\right)^{*} \cap\left(R_{N}^{1}\right)_{k+2 j-\frac{1}{2}}^{\text {trunc }} \\
& \text { and } \left.Z_{N}^{l-1, l} \stackrel{*}{\leftrightarrow} Z_{N}^{l, l+1} \text { in } T_{N+\frac{1}{2}, k+j+\frac{1}{2}}^{l} \backslash \mathscr{B}\left(Z_{N}\right) \text { for } l=2,3,4\right) \\
& \leqslant c_{25}^{2 \epsilon_{0} j} \mathbb{P}_{N, k}^{0}\left(\hat{C}_{-}\left(\hat{\Gamma}_{1,2}, \sigma^{x}\right)=J \text {, and in } \omega^{x}, \psi_{1,2}-\frac{1}{2} e_{1} \stackrel{*}{\leftrightarrows} Z_{N}^{4,1}\right. \\
& \text { in } \overline{\mathscr{B}}\left(\Lambda_{N} \backslash \bar{J}\right)^{*} \cap\left(R_{N}^{1}\right)_{k+2 j-\frac{1}{2}}^{\text {trunc }} \\
& \text { and } Z_{N}^{l-1, l} \stackrel{*}{\leftrightarrow} Z_{N}^{l, l+1} \text { in } T_{N+\frac{1}{2}, k+j+\frac{1}{2}}^{l} \backslash \mathscr{B}\left(Z_{N}\right) \text { for } l=2,3,4 \\
& \text { and } \sigma_{y}=1 \text { for all } y \in \psi_{1,2} \cap \partial J \text { ) } \\
& =c_{25}^{2 \epsilon_{10} j} \mathbb{P}_{N, k}^{0}\left(\mathrm{in} \omega^{x}, \psi_{1,2}-\frac{1}{2} e_{1} \stackrel{*}{\stackrel{ }{*}} Z_{N}^{4,1} \text { in } \overline{\mathscr{B}}\left(\Lambda_{N} \backslash \bar{J}\right)^{*} \cap\left(R_{N}^{1}\right)_{k+2 j-\frac{1}{2}}^{\text {trunc }}\right. \\
& \text { and } Z_{N}^{l-1, l} \stackrel{*}{\leftrightarrow} Z_{N}^{l, l+1} \text { in } T_{N+\frac{1}{2}, k+j+\frac{1}{2}}^{l} \mathscr{B}\left(Z_{N}\right) \text { for } l=2,3,4 \mid \hat{C}_{-}\left(\hat{\Gamma}_{1,2}, \sigma^{x}\right)=J \\
& \text { and } \sigma_{y}=1 \text { for all } y \in \psi_{12} \cap \partial J \text { ) } \\
& \cdot \mathbb{P}_{N, k}^{0}\left(\hat{C}_{-}\left(\hat{\Gamma}_{1,2}, \sigma^{x}\right)=J \text { and } \sigma_{y}=1 \text { for all } y \in \psi_{1,2} \cap \partial J\right) .(3.20)
\end{aligned}
$$

Here $\psi_{1,2}-\frac{1}{2} e_{1}$ means the translate of $\psi_{1,2}$ by $-\frac{1}{2} e_{1}$. When $\hat{C}_{-}\left(\hat{\Gamma}_{1,2}, \sigma^{x}\right)=J$, we have by definition $\sigma_{y}^{x}=1$ for all $y \in \partial J \cap\left(\Lambda_{N}\right)_{k+2 j}^{\text {trunc }}$ and for all $y \in \hat{\Gamma}_{1,2} \backslash J$, so from the Markov property of the Ising model, the conditioning on the right side of (3.20) is equivalent to conditioning on $\sigma_{y}^{x}=1$ for all $y \in \bar{J} \cup \hat{\Gamma}_{1,2}$. Therefore the first probability on the right side of (3.20) is

$$
\begin{align*}
& P_{\Lambda_{N} \backslash\left(\bar{J} \cup \hat{I}_{1,2)}, \eta^{k, 0, J}\right.}^{p, 2}\left(\text { in } \omega^{x}, \psi_{1,2}-\frac{1}{2} e_{1} \stackrel{*}{\leftrightarrow} Z_{N}^{4,1} \text { in } \overline{\mathscr{B}}\left(\Lambda_{N} \backslash\left(\bar{J} \cup \hat{\Gamma}_{1,2}\right)\right)^{*} \cap\left(R_{N}^{1}\right)_{k+2 j-\frac{1}{2}}^{\text {trunc }}\right. \\
& \left.\quad \text { and } Z_{N}^{l-1, l} \stackrel{*}{\leftrightarrow} Z_{N}^{l, J+1} \text { in } T_{N+\frac{1}{2}, k+j+\frac{1}{2}}^{l} \backslash \mathscr{B}\left(Z_{N}\right) \text { for } l=2,3,4\right) . \tag{3.21}
\end{align*}
$$

Let

$$
\begin{aligned}
\hat{Z}_{N}^{1,2} & =\left\{x \in \mathbb{R}^{2}: d\left(x, \Gamma_{1,2} \backslash \hat{\Gamma}_{1,2}\right) \leqslant c_{21} \log N\right\}, \\
\hat{Z}_{N} & =\hat{Z}_{N}^{1,2} \cup Z_{N}^{2,3} \cup Z_{N}^{3,4} \cup Z_{N}^{4,1},
\end{aligned}
$$

where $c_{21}$ is from the definition of $Z_{N}$. One effect of changing the measure from $P_{\Lambda_{N}, \eta^{k}, 0}^{p, 2}$ to $P_{\Lambda_{N} \backslash\left(\bar{J} \cup \hat{r}_{1,2}\right), \eta^{k, 0, J}}^{p, 2}$ is to shrink the free-boundary influence region
from $Z_{N}$ to $\hat{Z}_{N}$. More precisely, as in (3.16), we have using Lemma 2.2 that (3.21) is at most

$$
\begin{align*}
& P_{A_{N} \backslash\left(\bar{J} \cup \hat{\Gamma}_{1,2}\right), \eta^{k, 0,},}^{p, ~(i n} \omega^{x}, Z_{N}^{4,1} \stackrel{*}{\leftrightarrow}\left(\psi_{1,2}-\frac{1}{2} e_{1}\right) \cup \hat{Z}_{N}^{1,2} \\
& \text { in } \overline{\mathscr{B}}\left(\Lambda_{N} \backslash\left(\bar{J} \cup \hat{\Gamma}_{1,2}\right)\right)^{*} \cap\left(R_{N}^{1}\right)_{k+2 j-\frac{1}{2}}^{\text {trunc }} \backslash \mathscr{B}\left(\hat{Z}_{N}\right) \text { and } \\
& \left.Z_{N}^{l-1, l} \stackrel{*}{\leftrightarrow} Z_{N}^{l, l+1} \text { in } T_{N+\frac{1}{2}, k+j+\frac{1}{2}}^{l} \mathscr{B}\left(\hat{Z}_{N}\right) \text { for } l=2,3,4\right) \\
& \leqslant 16 P^{p, 2}\left(\text { in } \omega^{x}, Z_{N}^{4,1} \stackrel{*}{\leftrightarrows}\left(\psi_{1,2}-\frac{1}{2} e_{1}\right) \cup \hat{Z}_{N}^{1,2}\right. \\
& \text { in } \left.\overline{\mathscr{B}}\left(\Lambda_{N} \backslash\left(\bar{J} \cup \hat{\Gamma}_{1,2}\right)\right)^{*} \cap\left(R_{N}^{1}\right)_{k+2 j-\frac{1}{2}}^{\text {trunc }} \backslash \mathscr{B}\left(\hat{Z}_{N}\right)\right) \\
& \cdot \prod_{l=2}^{4} P^{p, 2}\left(Z_{N}^{l-1, l} \stackrel{*}{\leftrightarrow} Z_{N}^{l, l+1} \text { in } T_{N+\frac{1}{2}, k+j+\frac{1}{2}}^{l} \backslash \mathscr{B}\left(\hat{Z}_{N}\right)\right) . \tag{3.22}
\end{align*}
$$

To bound the first probability on the right side of (3.22), we observe that, provided $K$ is large, if $u \in \hat{Z}_{N}^{4,1}$ and $v \in\left(\psi_{1,2}-\frac{1}{2} e_{1}\right) \cup \hat{Z}_{N}^{1,2}$, then

$$
\tau(v-u) \geqslant\left(2 k+2 j-\frac{1}{2}-2 c_{21} \log N\right) \tau\left(e_{1}\right) \geqslant(2 k+j) \tau\left(e_{1}\right) .
$$

Further, we can replace $\omega^{x}$ with $\omega$ at the expense of at most a constant factor. Therefore as in (3.16), the right side of (3.22) is at most

$$
c_{26} N^{2} e^{-(2 k+j) \tau\left(e_{1}\right)}\left(N^{c_{23}} e^{-2 k \tau\left(e_{1}\right)}\right)^{3} \leqslant e^{-\left(8 k+\frac{1}{2} j\right) \tau\left(e_{1}\right)} .
$$

Plugging this into (3.20), provided $\epsilon_{10}<1 / 8$ and $K$ is large we obtain

$$
\begin{align*}
\mathbb{P}_{N, k}^{0}\left(A_{3}^{\prime}\right) & \leqslant \sum_{x, z, J} c_{25}^{2 \epsilon_{15} j} e^{-\left(8 k+\frac{1}{2} j\right) \tau\left(e_{1}\right)} \mathbb{P}_{N, k}^{0}\left(\hat{C}_{-}\left(\hat{\Gamma}_{1,2}, \sigma^{x}\right)=J\right) \\
& \leqslant e^{-\left(8 k+\frac{1}{4} j\right) \tau\left(e_{1}\right)} . \tag{3.23}
\end{align*}
$$

A similar proof gives the same bound for $\mathbb{P}_{N, k}^{0}\left(A_{3}^{\prime \prime}\right)$. Combining this with (3.16) and (3.17) gives

$$
\begin{equation*}
\left.\mu_{A_{N}, \eta^{k}, 0}^{\beta} \partial_{i n} D\right)=\mathbb{P}_{N, k}^{0}\left(A_{1} \cup A_{2} \cup A_{3}\right) \leqslant e^{-\left(8 k+\epsilon_{11} j\right) \tau\left(e_{1}\right)} . \tag{3.24}
\end{equation*}
$$

It follows easily from (3.14) that $\mu_{A_{N}, \eta^{k, o}}^{\beta}(D) \leqslant 1 / 2$. Combining this with (3.12), (3.24) and (1.2) yields

$$
\Delta\left(\Lambda_{N}, \eta^{k, 0}, \beta\right) \leqslant c_{27}=\left|\Lambda_{N}\right| k^{4 c_{13}} e^{-\epsilon_{11} j} \leqslant e^{-\epsilon_{11} j / 2}=e^{-\epsilon_{11}(N-k) / 8},
$$

which proves Theorem 1.1 for $\epsilon=0$.

## 4. PROOF OF THEOREMS 1.1(ii) AND 1.2

The FK measure corresponding to the boundary condition $\eta^{k,-1}$ is given (cf. (2.2)) by

$$
P_{A_{N}, \eta^{k,-1}}^{p, 2}=P_{\Lambda_{N}, w}^{p, 2}\left(\cdot \mid V\left(\Lambda_{N}, \eta^{k,-1}\right)\right),
$$

where $V\left(\Lambda_{N}, \eta^{k,-1}\right)$ is the event that there is no open dual path from $\Gamma_{i}$ to $\Gamma_{j, j+1}$ for any $i \neq j$. Our calculations, however, are facilitated by using a different conditioning, as follows. Consider bond configurations on $\bar{B}\left(\Lambda_{N+1}\right)$. Let $U^{\mathrm{FK}}$ denote the event that for all $i$, all bonds $\langle x y\rangle$ with $x, y \in \Gamma_{i}$ are open, all bonds $\langle x y\rangle$ with $x, y \in \Gamma_{i, i+1}$ are open, and all other bonds in $\overline{\mathscr{B}}\left(\Lambda_{N+1}\right) \backslash \overline{\mathscr{B}}\left(\Lambda_{N}\right)$ are closed. The FK model $P_{A_{N+1}, f}^{p, 2}\left(\cdot \mid U^{\mathrm{FK}}\right)$ corresponds to an Ising model $\mu_{\Lambda_{N+1}, f}^{\beta}\left(\cdot \mid U^{\text {Ising }}\right)$, where $U^{\text {Ising }}$ is the event that for all $i$, all sites in $\Gamma_{i}$ have the same spin, and all sites in $\Gamma_{i, i+1}$ have the same spin. Let $L$ denote the event that for all $i$, all sites in $\Gamma_{i}$ have spin 1 and all sites in $\Gamma_{i, i+1}$ have spin -1 . Then

$$
\mu_{\Lambda_{N+1}, f}^{\beta}\left(\sigma_{\Lambda_{N}} \in \cdot \mid L\right)=\mu_{\Lambda_{N}, \eta^{k,-1}} .
$$

The measure $\mathbb{P}_{N+1, k}^{0}\left(\cdot \mid \omega \in U^{\mathrm{FK}}\right)=\mathbb{P}_{N+1, k}^{0}\left(\cdot \mid \sigma \in U^{\text {Ising }}\right)$ gives the joint construction, coupling $P_{A_{N+1, k}, f}^{p_{2}}\left(\cdot \mid U^{\mathrm{FK}}\right)$ and $\mu_{A_{N+1}, f}^{\beta}\left(\cdot \mid U^{\text {Ising }}\right)$.

For $\sigma \in L$, let

$$
\begin{aligned}
\mathscr{J} & =\{(i, j): 1 \leqslant i, j \leqslant 4, i<j\}, \\
\mathscr{A}_{N}(\sigma) & =\left\{(i, j) \in \mathscr{J}: \Gamma_{i} \stackrel{+}{\leftrightarrows} \Gamma_{j} \text { in } \overline{\mathscr{B}}\left(\Lambda_{N}\right)\right\} .
\end{aligned}
$$

As motivation, note we expect that, roughly,

$$
\begin{array}{rll}
\mu_{A_{N+1}, f}^{\beta}\left(\mathscr{A}_{N}=\mathscr{J} \mid L\right) \approx 1 & \text { if } & 2 k \tau\left(e_{1}\right)>(N-k) \tau\left(e_{1}+e_{2}\right),  \tag{4.1}\\
\mu_{A_{N+1}, f}^{\beta}\left(\mathscr{A}_{N}=\phi \mid L\right) \approx 1 & \text { if } & 2 k \tau\left(e_{1}\right)<(N-k) \tau\left(e_{1}+e_{2}\right) .
\end{array}
$$

In the case $2 k \tau\left(e_{1}\right) \geqslant(N-k) \tau\left(e_{1}+e_{2}\right)$, we will bound the spectral gap using in (1.2) the same event $D$ as in Section 3, but in the opposite case we replace it with a different event $\hat{D}=\bigcap_{i=1}^{4} \hat{D}_{i, i+1}$. Here $\hat{D}_{1,2}$ is the event that there is no minus-path in $\sigma$ in $\overline{\mathscr{B}}\left(\Lambda_{N}\right)$ from $\Gamma_{1,2}$ to $\left(S_{N, k / 4}^{1,2}\right)^{c}$, where $S_{N, m}^{1,2}$ is the square $[m, N+1] \times[-N-1,-m]$, and $\hat{D}_{i, i+1}, S_{N, m}^{i, 1}$ are the corresponding event and square obtained by rotation, for $i=2,3,4$.

Suppose first that $2 k \tau\left(e_{1}\right) \geqslant(N-k) \tau\left(e_{1}+e_{2}\right)$. Let $x_{1,1}=\left(-k-\frac{1}{2}\right.$, $\left.-N-\frac{1}{2}\right)$ and $x_{1,2}=\left(k+\frac{1}{2},-N-\frac{1}{2}\right)$. These dual sites are approximately at
the ends of $\Gamma_{1}$. We define corresponding sites $x_{i j}$ for $i=2,3,4$ and $j=1,2$. In place of the event $E_{i}$ of Lemma 3.1, we will use

$$
\hat{E}_{i}=\left\{\omega: x_{1,1} \stackrel{*}{\mapsto} x_{1,2} \text { in } T_{N+\frac{1}{2}, k+j+\frac{1}{2}}^{i}\right\} .
$$

Lemma 3.1 and its proof remain valid for $\hat{E}_{i}$ in place of $E_{i}$, and the proof of the lower bound (3.12) for $\mu_{\Lambda_{N}, \eta^{k}, 0}^{\beta}(D)$ goes through with minimal changes to give

$$
\mu_{\Lambda_{N+1}, f}^{\beta}\left(D \cap L \mid U^{\text {Ising }}\right) \geqslant \frac{\epsilon_{12}}{k^{c_{28}}} e^{-8 k \tau\left(e_{1}\right)} .
$$

The proof of (3.24) also goes through with minimal changes; in fact Case 3 can be made simpler using the fact that the boundary regions $\Gamma_{i, i+1}$ are each wired. (We will not do so here, since it is unnecessary.) The result is that

$$
\begin{equation*}
\mu_{\Lambda_{N+1}, f}^{\beta}\left(\partial_{i n} D \cap L \mid U^{\text {Ising }}\right) \leqslant e^{-\left(8 k+\epsilon_{13} j\right) \tau\left(e_{1}\right)} . \tag{4.3}
\end{equation*}
$$

Combining (4.2) and (4.3) gives

$$
\begin{align*}
\frac{\mu_{\Lambda_{N}, \eta^{k},-1}^{\beta}\left(\partial_{i n} D\right)}{\mu_{A_{N}, \eta^{k},-1}^{\beta}(D)} & =\frac{\mu_{\Lambda_{N+1}, f}^{\beta}\left(\partial_{i n} D \mid L\right)}{\mu_{A_{N+1}, f}^{\beta}(D \mid L)} \\
& =\frac{\mu_{\Lambda_{N+1}, f}^{\beta}\left(\partial_{i n} D \cap L \mid U^{\text {Ising }}\right)}{\mu_{\Lambda_{N+1}, f}^{\beta}\left(D \cap L \mid U^{\text {Ising }}\right)} \\
& \leqslant \frac{\epsilon_{12}}{k^{\alpha_{28}}} e^{-\epsilon_{13} j \tau\left(e_{1}\right)} . \tag{4.4}
\end{align*}
$$

In Section 3 we easily obtained the lower bound $\mu_{A_{N}, \eta^{k}, 0}^{\beta}\left(D^{c}\right) \geqslant 1 / 2$ to complete the proof. Here the situation is a little more complex. A lower bound of the form

$$
\begin{equation*}
\frac{\mu_{\Lambda_{N+1}, f}^{\beta}\left(D^{c} \cap L \mid U^{\text {Ising }}\right)}{\mu_{A_{N+1}, f}^{\beta}\left(D \cap L \mid U^{\text {Ising }}\right)} \geqslant \theta \tag{4.5}
\end{equation*}
$$

for some $\theta$ is equivalent to the statement

$$
\begin{equation*}
\mu_{\Lambda_{N}, \eta^{k},-1}^{\beta}\left(D^{c}\right)=\mu_{\Lambda_{N+1}, f}^{\beta}\left(D^{c} \mid L\right) \geqslant \frac{\theta}{1+\theta} . \tag{4.6}
\end{equation*}
$$

Hence we consider bounds for the numerator and denominator of (4.5). Let $F_{i, i+1}$ denote the event that $x_{i, 2} \stackrel{*}{\leftrightarrow} x_{i+1,1}$ via a path in $S_{N, k / 4}^{i, i+1}$, and $F=\bigcap_{i=1}^{4} F_{i, i+1}$. We have

$$
\begin{align*}
\mu_{A_{N+1}, f}^{\beta}\left(D^{c} \cap L \mid U^{\mathrm{Ising}}\right) & \geqslant \mu_{\Lambda_{N+1}, f}^{\beta}\left(\hat{D} \cap L \mid U^{\mathrm{Ising}}\right) \\
& \geqslant P_{A_{N+1}, f}^{p, 2}\left(F \mid U^{\mathrm{FK}}\right) \mathbb{P}_{N+1, k}^{0}\left(\hat{D} \cap L \mid F \cap U^{\mathrm{FK}}\right) . \tag{4.7}
\end{align*}
$$

It is straightforward to prove an analog of Lemma 2.5 for $S_{N+\frac{1}{2}, 2 m+\frac{1}{2}}^{1,2}$ in place of $T_{N+\frac{1}{2}, 2 m+\frac{1}{2}}^{1}$. Therefore mimicking the proof of Lemma 3.1, we obtain

$$
\begin{equation*}
P_{\Lambda_{N+1}, f}^{p, 2}\left(F_{i, i+1} \mid U^{\mathrm{FK}}\right) \geqslant \frac{\epsilon_{14}}{(N-k)^{c_{29}}} e^{-(N-k) \tau\left(e_{1}+e_{2}\right)} . \tag{4.8}
\end{equation*}
$$

Then, analogously to (3.12), from (4.7),

$$
\begin{equation*}
\mu_{\Lambda_{N+1}, f}^{\beta}\left(D^{c} \cap L \mid U^{\text {Ising }}\right) \geqslant \frac{\epsilon_{15}}{(N-k)^{c_{29}}} e^{-4(N-k) \tau\left(e_{1}+e_{2}\right)} . \tag{4.9}
\end{equation*}
$$

Next we have, using (2.3) and Lemma 2.2,

$$
\begin{align*}
\mu_{\Lambda_{N+1}, f}^{\beta} & \left(D \cap L \mid U^{\text {Ising }}\right) \\
& \leqslant P_{A_{N+1}, f}^{p, 2}\left(x_{i, 2} \stackrel{*}{\mapsto} x_{i+1,1} \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} \text { for all } i \mid U^{\mathrm{FK}}\right) \\
& \leqslant 2^{8} P_{\Lambda_{N}, w}^{p, 2}\left(x_{i, 2} \stackrel{*}{\leftrightarrow} x_{i+1,1} \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} \text { for all } i\right) \\
& \leqslant 2^{8} P^{p, 2}\left(x_{i, 2} \stackrel{*}{\leftrightarrow} x_{i+1,1} \text { in } T_{N+\frac{1}{2}, k+3 j+\frac{3}{2}}^{i} \text { for all } i\right) \\
& \leqslant 2048 e^{-8 k t\left(e_{1}\right)} . \tag{4.10}
\end{align*}
$$

Since $2 k \tau\left(e_{1}\right) \geqslant(N-k) \tau\left(e_{1}+e_{2}\right)$, (4.9) and (4.10) give

$$
\frac{\mu_{\Lambda_{N+1}, f}^{\beta}\left(D^{c} \cap L \mid U^{\text {Ising }}\right)}{\mu_{\Lambda_{N+1}, f}^{\beta}\left(D \cap L \mid U^{\text {Ising }}\right)} \geqslant \frac{\epsilon_{15}}{8(N-k)^{c_{29}}} .
$$

With (4.5) and (4.6), this shows

$$
\mu_{\Lambda_{N}, \eta^{k,-1}}^{\beta}\left(D^{c}\right) \geqslant \frac{\epsilon_{15}}{16(N-k)^{c_{29}}},
$$

which with (4.4) completes the proof of Theorem 1.1 for $\epsilon=-1$, as in Section 3.

The proof when $2 k \tau\left(e_{1}\right)<(N-k) \tau\left(e_{1}+e_{2}\right)$ is similar, with the roles of $D$ and $\hat{D}$ interchanged, using squares $S^{i, i+1}$ in place of the triangles $T^{i}$,.

Turning to Theorem 1.2, Let us write $\mu$ for $\mu_{1, \eta}^{\beta}$ and $\mu^{\prime}$ for $\mu_{1, \eta^{\prime}}^{\beta}$. There exists a constant $M=M(\beta)$ such that for every $\sigma \in \Sigma_{A}$, we have $\mu^{\prime}(\sigma) \leqslant M \mu(\sigma)$. Therefore, for every nonnegative function $g$ on $\Sigma_{A}$,

$$
\langle g\rangle_{\mu^{\prime}} \leqslant M\langle g\rangle_{\mu} .
$$

Hence for $f \in L^{2}\left(\mu^{\prime}\right)=L^{2}(\mu)$,

$$
\operatorname{var}_{\mu^{\prime}}(f)=\left\langle\left(f-\langle f\rangle_{\mu^{\prime}}\right)^{2}\right\rangle_{\mu^{\prime}} \leqslant\left\langle\left(f-\langle f\rangle_{\mu}\right)^{2}\right\rangle_{\mu^{\prime}} \leqslant M \operatorname{var}_{\mu}(f) .
$$

Interchanging the roles of $\mu$ and $\mu^{\prime}$ we see that the corresponding Dirichlet forms $\mathscr{D}_{\mu}$ and $\mathscr{D}_{\mu^{\prime}}$ satisfy

$$
-\mathscr{D}_{\mu^{\prime}}(f, f) \geqslant-M \mathscr{D}_{\mu}(f, f) .
$$

With (1.1) this proves (i), and (ii) is a straightforward consequence of (i).

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